

Vector-boson pair production at the LHC

Tobias Kasprzik

In collaboration with A. Bierweiler, S. Gieseke, J. H. Kühn

Karlsruhe Institute of Technology (KIT),
Institut für Theoretische Teilchenphysik (TTP)

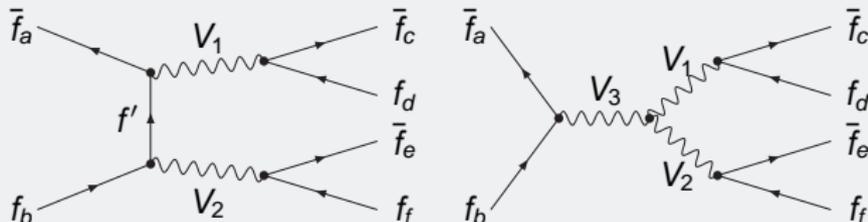
[[arXiv:1208.3147](https://arxiv.org/abs/1208.3147), [arXiv:1208.3404](https://arxiv.org/abs/1208.3404), [arXiv:1305.5402](https://arxiv.org/abs/1305.5402)]

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- 1 Introduction
- 2 Status of Theory Predictions
- 3 NLO EW Corrections
- 4 Electroweak Corrections in HERWIG++
- 5 Summary & Outlook

Vector-boson pair production: $pp \rightarrow WW/ZZ/WZ \rightarrow 4l$

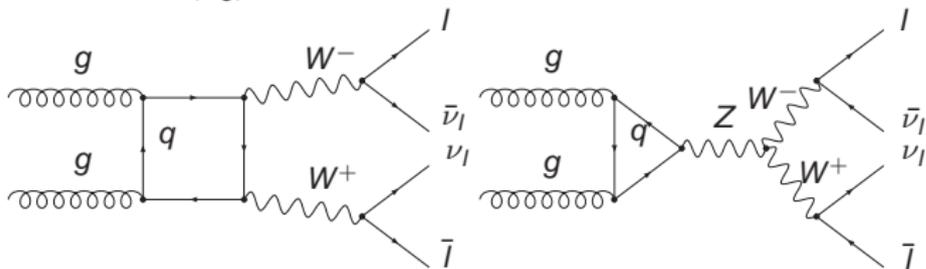


- ZZ/WW/ $\gamma\gamma$ production important **irreducible background** to inclusive SM Higgs-boson production
- Probe **non-abelian structure** of the Standard Model (SM) at high energies
- Search for **anomalous couplings**
- Backgrounds to **new-physics searches**, i.e. leptons + \cancel{E}_T signatures
→ SUSY-particle pair production

Extensive study of production of $WW, WZ, ZZ, W\gamma, Z\gamma, \gamma\gamma$ at NLO QCD [Campbell, Ellis, Williams

'05; Campbell, Ellis '99; Dixon, Kunszt, Signer '98]

- Results matched with **parton showers** \oplus combined with **soft gluon resummation**
[Nason, Ridolfi '06; Frixione, Webber '06; Grazzini '06; Dawson, Lewis, Zeng '13]
- On-shell **leptonic decays** of the vector bosons taken into account (narrow-width approximation) retaining all spin information
- Cross section at LO dominated by the $q\bar{q}$ channels
 - Significant contributions of the channels $gg \rightarrow V_1 V_2 \sim 10\%$ to LO $q\bar{q}$, although formally at $\mathcal{O}(\alpha_s^2)$ [Glover, van der Bij '89; Kao, Dicus '91; Duhrssen et al. '05]

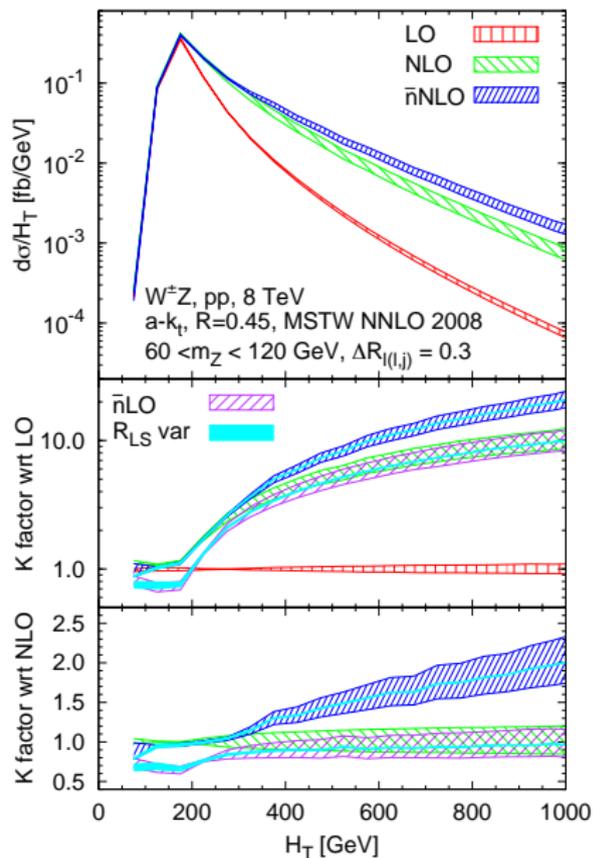


- **Even larger corrections of 30% if event selection for Higgs searches is applied** [Binoth et al. '06]

Huge NLO K -factors at high p_T in V -pair production, large residual uncertainties

→ need NNLO for accurate predictions!

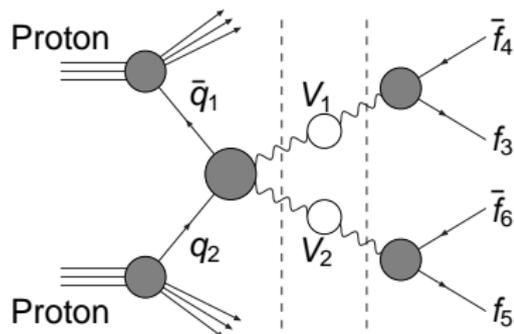
- ✓ NNLO QCD for $pp \rightarrow \gamma\gamma/Z\gamma$ **known fully differentially** [Bern, de Freitas, Dixon '01; Catani, Cieri, de Florian, Ferrera, Grazzini '12; Grazzini, Kallweit, Rathlev, Torre '13]
- **Two-loop matrix** elements known for $V\gamma$ [Gehrmann, Tancredi, Weihs, '12/13] and WW (high-energy approx.) [Chachamis, Czakon, Eiras '08]
Recently: Two-loop master integrals for $q\bar{q} \rightarrow VV$ (planar topologies) [Gehrmann, Tancredi, Weihs '13]
- $pp \rightarrow VV + \text{jet}$ known at NLO [Dittmaier, Kallweit, Uwer '08; Campanario, Englert, Spannowsky, Zeppenfeld '09/10/11, . . .]
- **Missing:**
 - non-planar topologies at 2 loop
 - double-real radiation with 2 soft/collinear partons
 - $pp \rightarrow VV + \text{jet}$ at one loop with one soft/collinear parton
- **Approximate NNLO result** for WZ/WW production provided recently [Campanario, Sapeta '12, Campanario, Rauch, Sapeta '13] using the LoopSim method [Rubin, Salam, Sapeta '10] (**caveat:** only reliable at high p_T)



[Campanario, Sapeta [arXiv:1209.4595v1]]

- Huge NLO K -factors at high p_T
- Prominent shift going from NLO to \bar{n} NLO (larger than NLO scale uncertainty!)
- Reduction of scale uncertainty
- Residual uncertainty due to missing NNLO terms small!
- Low p_T : method does not work, full NNLO still needed!

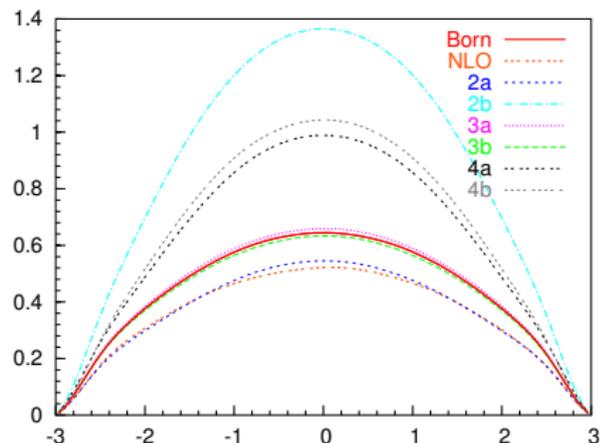
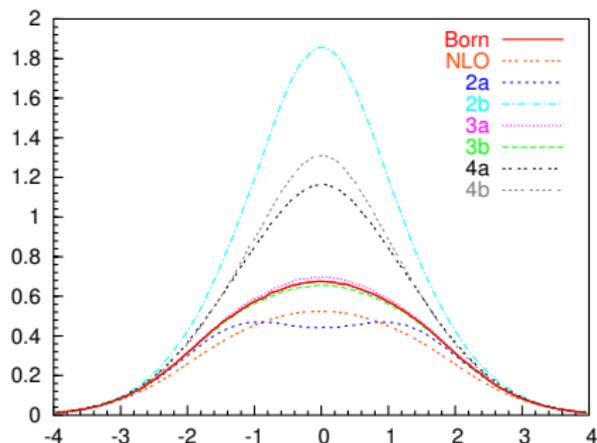
- $\mathcal{O}(\alpha)$ high-energy approximation known for all channels, vector bosons treated in **pole-approximation** → final-state leptons phenomenologically accessible [Accomando et al. '02-'06]



- **Full $\mathcal{O}(\alpha)$ corrections** known for $W\gamma$ and $Z\gamma$ production in single-pole approximation [Accomando, Denner, Maier '05]
- **NNLL effects** at two loops published for the WW channel [Kühn, Metzler, Uccirati, Penin '11]
- **We have calculated the full one-loop corrections to on-shell VV production at the LHC** [Bierweiler, TK, Kühn, Uccirati '12/13]
- Detailed NLO analysis of massive V-pair production [Baglio, Le Duc Ninh, Weber '13]

Consider anomalous WWZ coupling in $pp \rightarrow W^\pm Z \rightarrow l\bar{\nu}_l l' \bar{l}'$ at LHC14:
 ($p_{T,l} > 70$ GeV) [Accomando, Kaiser [arXiv:hep-ph/0511088]],

$\Delta y(Z, l)$ and $y(Z)$ distribution



- Significant distortion of distributions through aTGCs and EW corrections
- EW corrections may be misinterpreted as signal of aTGCs
 → EW corrections have to be included in aTGC analysis!

High-energy limit

$$s, |t|, |u| \gg M_V^2, \quad V = W, Z$$

→ **bosons have to be produced at large p_T**

- EW corrections at high energies dominated by **universal large logarithms**

$$\propto \alpha^L \ln^{2L}(M_V/\sqrt{s}) \quad (\text{LL}),$$

$$\propto \alpha^L \ln^{2L-1}(M_V/\sqrt{s}) \quad (\text{NLL}), \dots$$

at the L -loop level

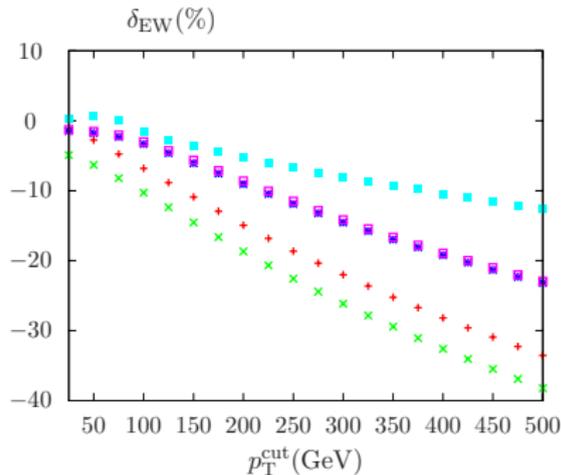
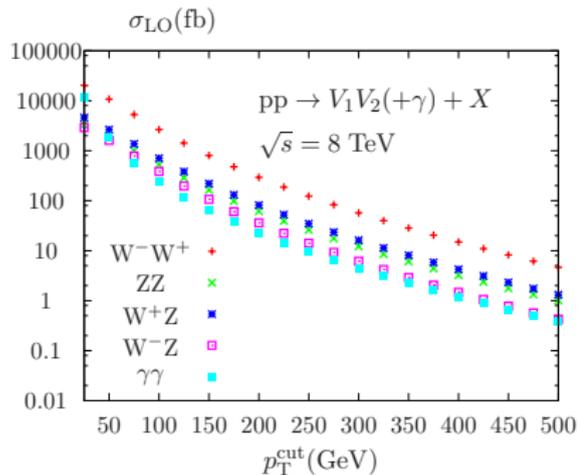
- **Corrections of $\sim -50\%$ at $\sqrt{s} = 2 \text{ TeV}$** (W-pair production)
- **Change of sign** going from LL to NLL (to NNLL . . .)
→ **substantial cancellations possible!**

- 1 Only consider **on-shell** vector bosons \oplus include all **mass effects**
- 2 Include **leptonic decays** \rightarrow **physical final states phenomenologically accessible**

Setup

- **Renormalization:**
On-shell scheme (G_μ , M_W , M_Z) to obtain UV finite
- **Virtual corrections:**
IR divergent (regularized by m_γ , m_q), compensated by
- **Real radiation:**
remaining collinear singularities to be absorbed in PDFs
- **Practical implementation:**
use MSTW2008LO PDFs [Martin et al. '09]
(impact of QED and factorization scheme small, in general sub-percent)

LHC at 8 TeV, default cuts: $p_{T,V} > 15$ GeV, $|y_V| < 2.5$



- **LO:** Drastically decreasing cross sections for large p_T
- **NLO:**
 - Full result, i.e. virtual, soft, hard, and collinear photons included
 - All mass effects included
 - Corrections largest for ZZ , smallest for $\gamma\gamma$
- Results published in [arXiv:1208.3147](https://arxiv.org/abs/1208.3147), [arXiv:1208.3404](https://arxiv.org/abs/1208.3404)

Low energies: Phase-space and perturbative suppression of $pp \rightarrow V_1 V_2 + (W/Z)$
⇒ contribution below 1%

High energies: **Logarithmic enhancement** of additional soft/collinear W- or Z-boson radiation

⇒ Investigation of $V_1 V_2 + W/Z$ production as background to V pairs at large p_T , M_{VV}

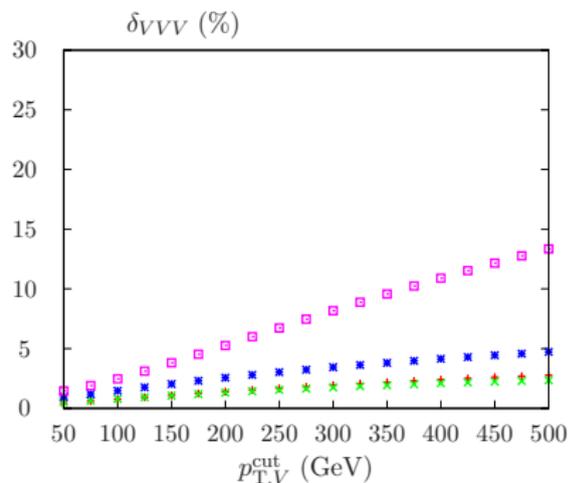
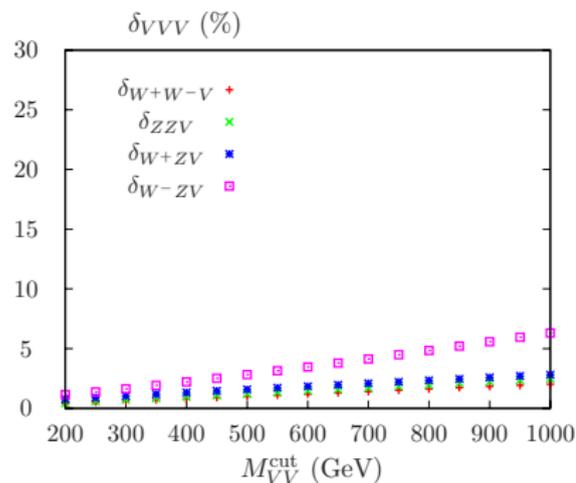
- invisible decay of $Z \rightarrow \nu\bar{\nu}$
- collinear emission
- ...

Simplified approach (details depend on experimental analysis),

e.g. W-Pair production:

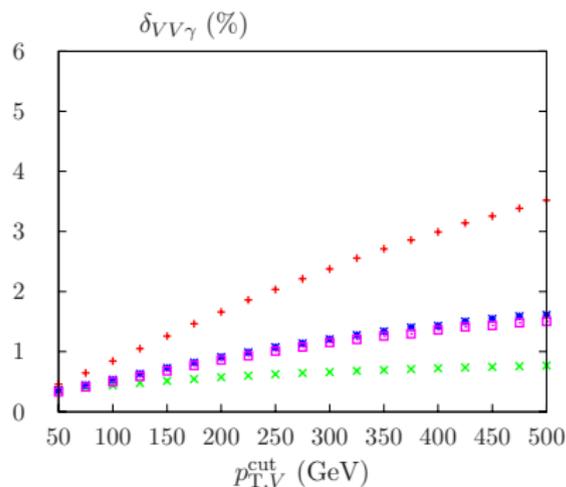
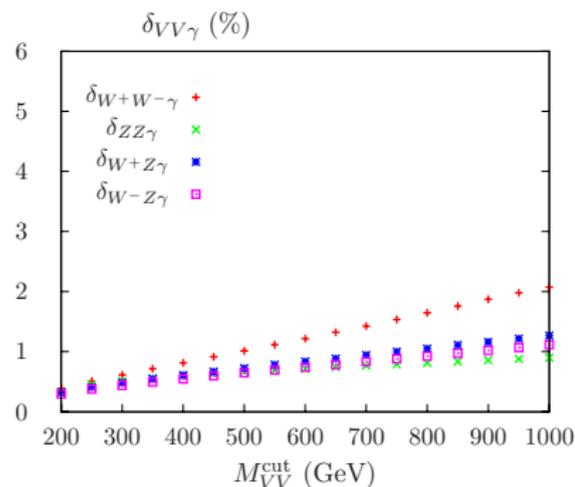
- 1 Include $pp \rightarrow W^- W^+ Z$ with totally inclusive Z
- 2 Include $pp \rightarrow W^- W^+ W^\pm$; treat W^\pm with lowest p_T totally inclusively

LHC at 8 TeV, default cuts: $p_{T,V} > 15$ GeV, $|y_V| < 2.5$



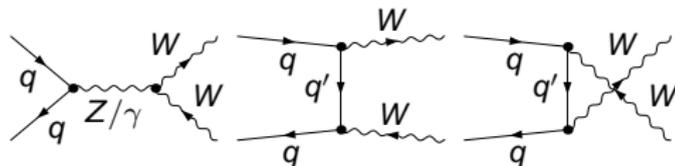
- corrections **below 5%** even for large transverse momenta and invariant masses
- corrections to W^-Z production enhanced due to W^-ZW^+ final states (PDFs!)

LHC at 8 TeV, default cuts: $p_{T,V} > 15$ GeV, $|y_V| < 2.5$, $p_{T,\gamma} > 15$ GeV, $|y_\gamma| < 2.5$

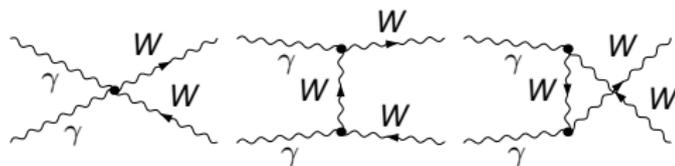


- real radiation of hard photons marginal ($< 2\%$)
→ neglect in MC implementation
- corrections largest for WW production

- **Partonic LO contributions at $\mathcal{O}(\alpha^2)$**

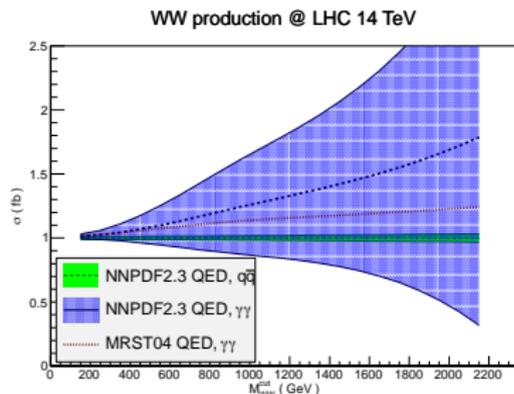
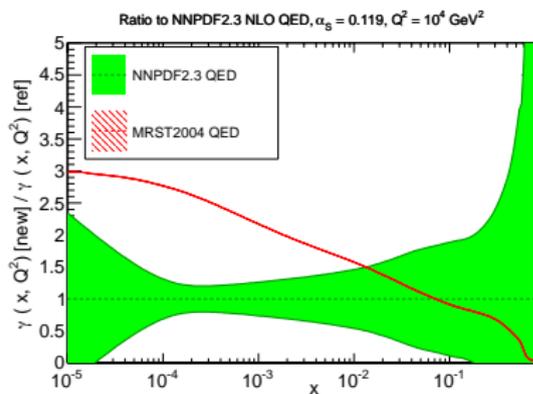


- **Photon-induced contributions at $\mathcal{O}(\alpha^2)$**



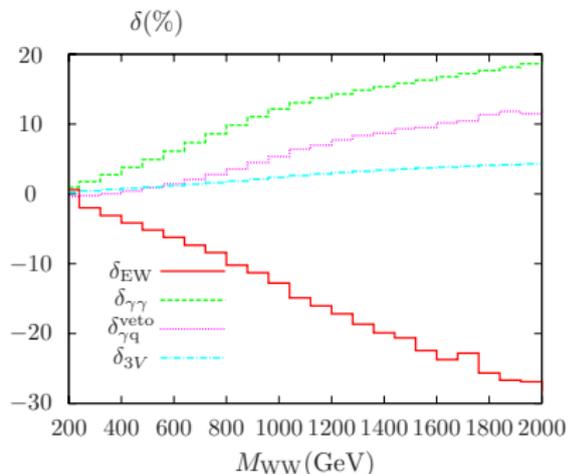
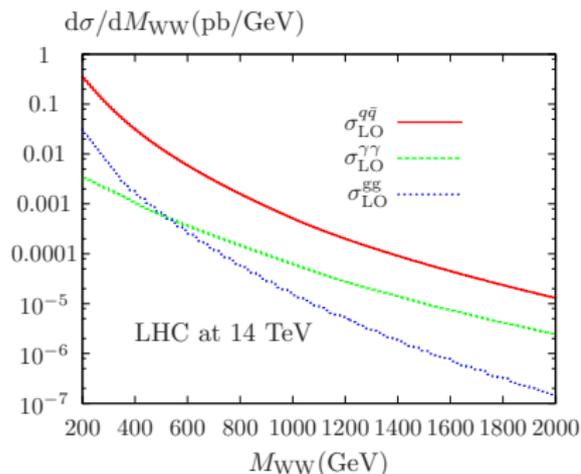
- Adopt MRST2004qed PDF set [Martin et al. '05]
 - ⊖ no error estimate
 - ⊖ no input from data for photon PDF
- Potentially large contribution at high invariant masses!
- **Now possible:** Comparison to **NNPDF2.3qed** [Carrazza '13], determined from DIS data

[Plots by Juan Rojo (thanks!)]



- Significant discrepancy between MRST2004_{qed} and NNPDF2.3_{qed} for photon PDFs
 - $\gamma\gamma \rightarrow WW$ at LHC14: Huge relative corrections at high invariant masses
 - Different predictions for MRST2004_{qed} (+20%) and NNPDF2.3_{qed} (+70%)
 - Huge error ($\sim \pm 50\%$) on NNPDF2.3_{qed} cross section
- further constrain photon PDFs through LHC WW production and DY [Carrazza '13 [arXiv: 1307.1131]]
- Potentially large effects from $q\gamma$ channels [Baglio, Le Duc Ninh, Weber '13]

LHC at 14 TeV, default cuts: $p_{T,V} > 15$ GeV, $|y_V| < 2.5$



- $\delta_{\gamma q}$ and $\delta_{\gamma\gamma}$ at the level of +10% and +20%
- δ_{EW} large and negative (−30%)

Compensation?

- **Photon-induced contributions:**

- $\sigma(\gamma\gamma \rightarrow WW) \rightarrow \frac{8\pi\alpha^2}{M_W^2}$ for large \hat{s}
 - ⇒ strong enhancement in forward & backward directions
- $q\gamma$ -induced contributions show flat angular distribution

- **weak corrections:**

negative Sudakov logs for large \hat{s} and \hat{t}

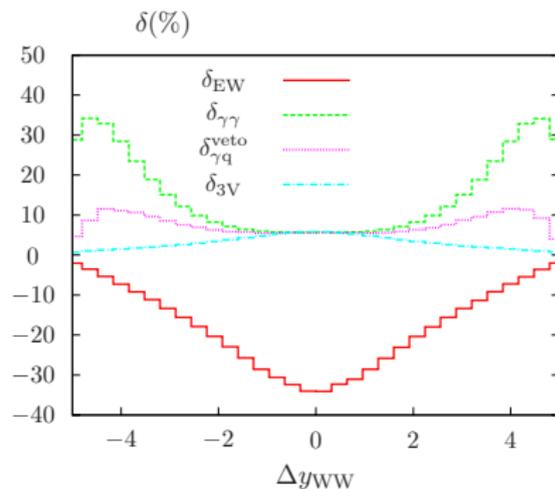
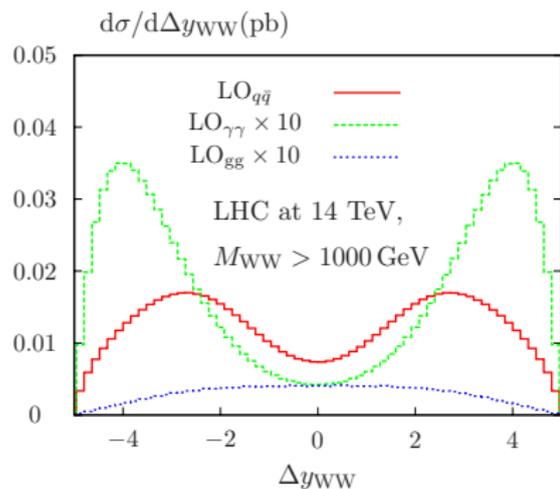
⇒ negative corrections for large scattering angles

- implications for $d\sigma/d\Delta y_{WW}$ with $\Delta y_{WW} = y_{W^+} - y_{W^-}$
(for fixed M_{WW} this corresponds to the angular distribution in the W-W rest frame!)

No compensation between photon-induced and weak corrections!

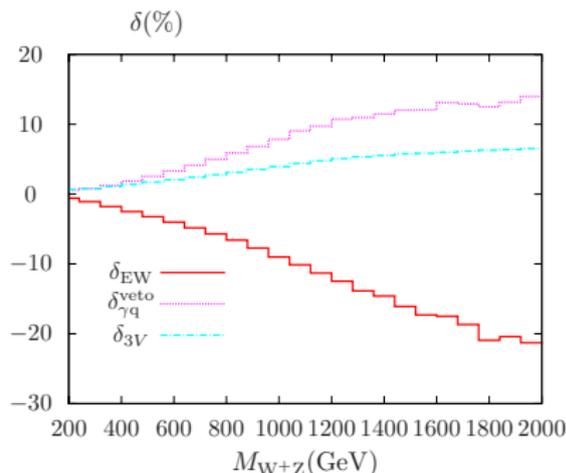
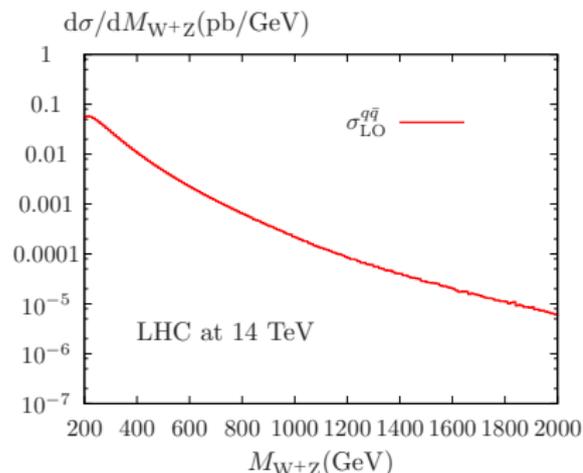
⇒ **Different angular distributions!**

W-pair production at high energies: $M_{WW} > 1000$ GeV



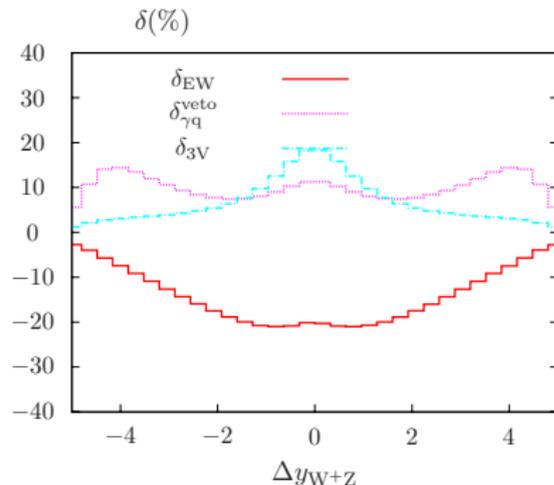
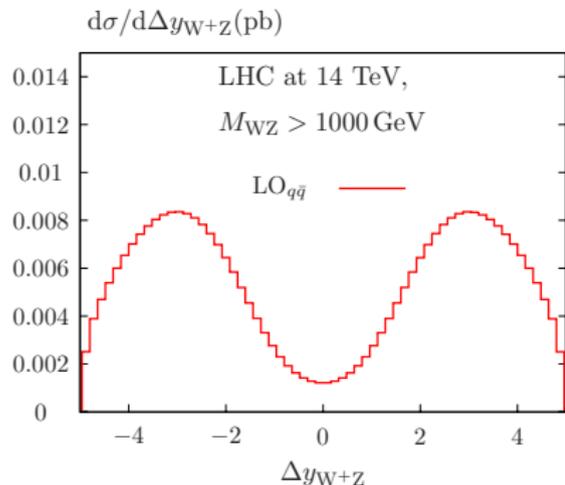
- strong forward-backward peaking of the angular distribution
- large photon-induced positive corrections at large scattering angles
- negative logarithmic weak corrections for central events
- drastic distortion of angular distributions → **anomalous couplings!**
- **Caveats:** $\delta_{\gamma q}$ and δ_{3V} depend on event selection.

LHC at 14 TeV, default cuts: $p_{T,V} > 15$ GeV, $|y_V| < 2.5$



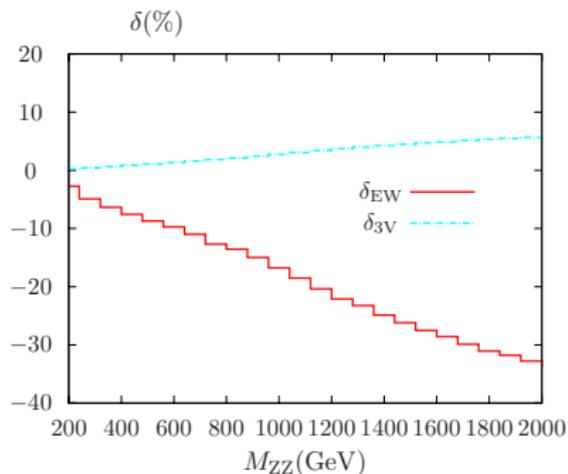
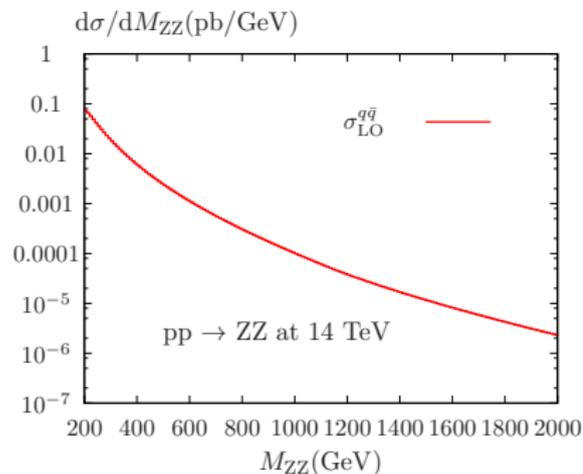
- moderate EW corrections at high invariant masses
- compensation of EW effects in invariant-mass spectrum
- **Caveats:** error on photon PDF, event selection in δ_{3V} and $\delta_{\gamma q}$, factorization of EW corrections . . .

W⁺Z production at high energies: $M_{WZ} > 1000$ GeV



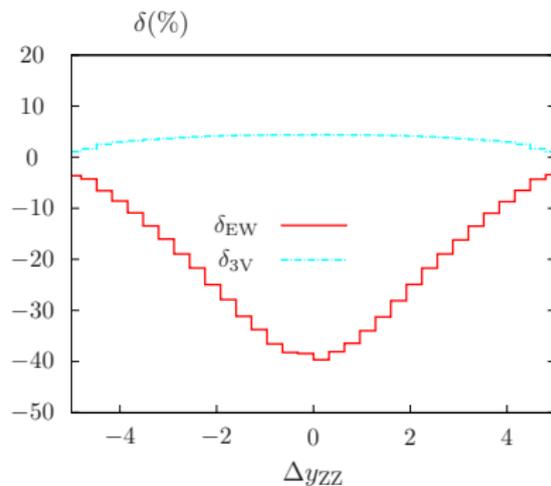
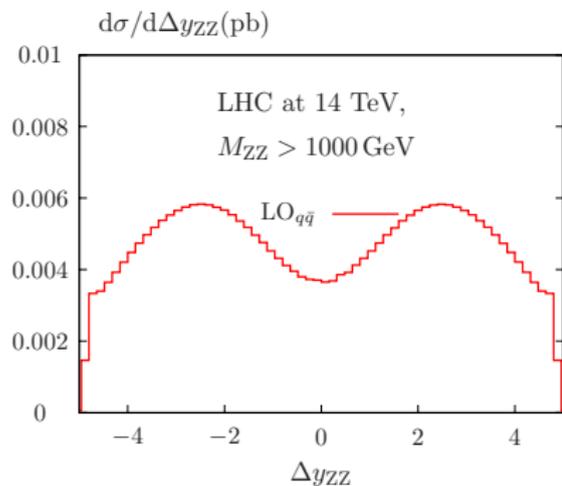
- drastic forward-backward peaking of the angular distribution
- partial compensation of photon-induced contributions, massive boson radiation and EW corrections

LHC at 14 TeV, default cuts: $p_{T,V} > 15$ GeV, $|y_V| < 2.5$



- large negative corrections at high invariant masses.
- photon-quark induced channels negligible.
- universal correction of -4% to the total cross section.

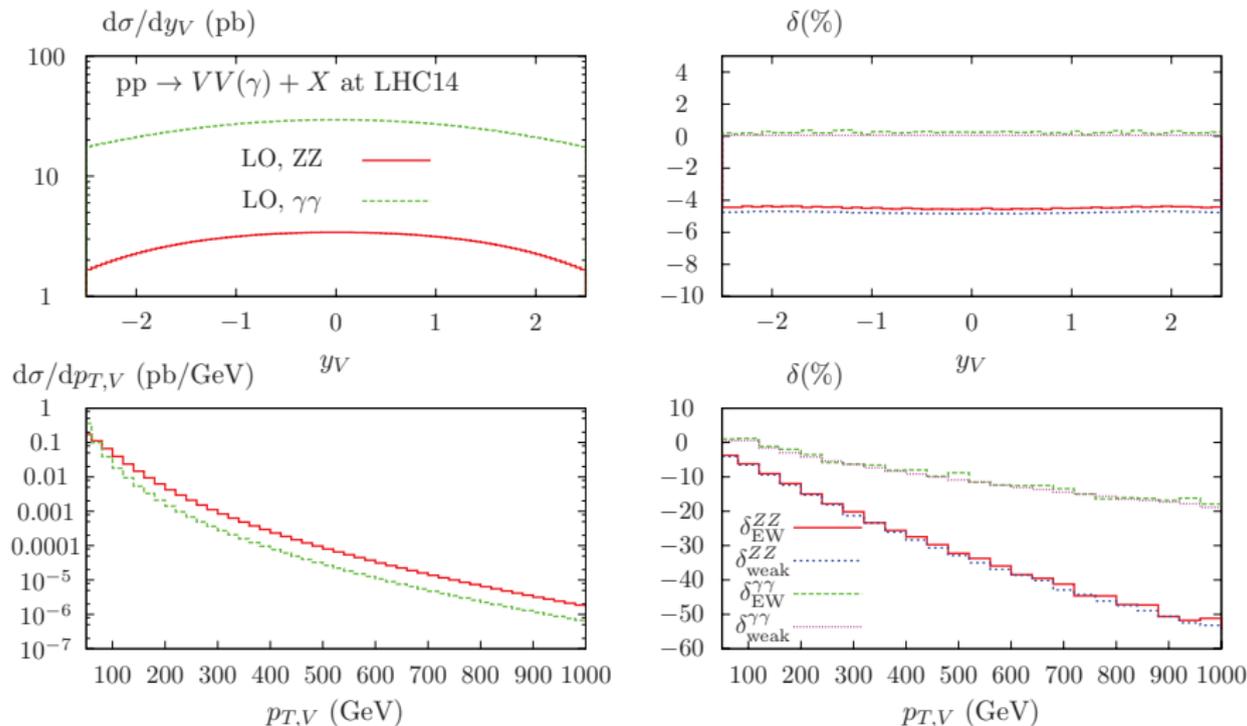
Z-pair production at high energies: $M_{ZZ} > 1000$ GeV



- drastic forward-backward peaking of the angular distribution
- prominent distortion of the angular distributions through EW corrections
- no pollution from massive boson radiation expected

Purely weak corrections well defined in ZZ/ $\gamma\gamma$ production

→ contributions of QED in general below 1%



Purely weak corrections well defined in ZZ/ $\gamma\gamma$ production
→ contributions of QED in general below 1%

- compute **purely weak corrections** to $pp \rightarrow (Z/\gamma^*)(Z/\gamma^*) \rightarrow e^+e^-\mu^+\mu^-$
- **LO: full calculation**, non-resonant and off-shell effects included.
 - naive fixed-width scheme
 - Complex-Mass Scheme (CMS) [Denner, Dittmaier, Roth, Wieders 2005]
- **NLO:** Two different approaches, including full spin correlations
 - **Double-Pole Approximation (DPA):** only doubly resonant contributions included, finite width taken into account
(On-shell projection, **caveat:** non-factorizable corrections neglected)
 - **Narrow-Width Approximation (NWA):** particles strictly on shell

$$\frac{1}{(Q^2 - M^2)^2 + M^2\Gamma^2} \rightarrow \frac{\pi}{M\Gamma} \delta(Q^2 - M^2),$$

valid if $\Gamma/M \rightarrow 0$.

$pp \rightarrow (Z/\gamma^*)(Z/\gamma^*) + X \rightarrow e^+e^-\mu^+\mu^- + X, \Delta y_{ZZ} < 3$					
$M_{inv}^{cut}(4l)/\text{GeV}$	$\sigma_{LO}^{naive}/\text{pb}$	$\sigma_{LO}^{CMS}/\text{pb}$	$\sigma_{LO}^{DPA}/\text{pb}$	$\sigma_{LO}^{NWA}/\text{pb}$	$\delta_{weak}^{DPA}/\%$
LHC14					
500	0.326×10^{-3}	0.326×10^{-3}	0.319×10^{-3}	0.343×10^{-3}	-15.9
600	0.168×10^{-3}	0.168×10^{-3}	0.164×10^{-3}	0.177×10^{-3}	-19.3
700	0.962×10^{-4}	0.962×10^{-4}	0.941×10^{-4}	1.017×10^{-4}	-22.3
800	0.587×10^{-4}	0.587×10^{-4}	0.575×10^{-4}	0.621×10^{-4}	-24.9
900	0.374×10^{-4}	0.374×10^{-4}	0.367×10^{-4}	0.397×10^{-4}	-27.4
1000	0.247×10^{-4}	0.247×10^{-4}	0.242×10^{-4}	0.262×10^{-4}	-29.7

LHC14, standard leptonic cuts

- **LO:** DPA works well, NWA: discrepancy of 5–10%
 - **NLO:** Good agreement ($\sim 1\%$) with K -factors obtained in Sudakov approximation [Accomando, Denner, Kaiser 2004]
- QED contributions (real-photon radiation, photon loops, non-factorizable contributions, corrections to Z-boson decay) **only at the 1% level**

Conclusion: weak K -factors of hard process sufficient to describe resonant 4-lepton production at reasonable accuracy

Our strategy

Factorization of EW and QCD corrections:

$$d\sigma_{\text{QCD}\times\text{EW}} = K_{\text{weak}}(\hat{s}, \hat{t}) \times d\sigma_{\text{QCD}}$$

σ_{QCD} : best prediction available for QCD-corrected cross section

- **Assumption:** bulk of EW effects properly described by weak K -factor $K_{\text{weak}}(\hat{s}, \hat{t})$ derived from $2 \rightarrow 2$ process.
- **FSR included** in YFS formalism (SOPHTY) [Hamilton, Richardson 2006] (only dressed leptons)
- $K_{\text{weak}}(\hat{s}, \hat{t})$ computed once and for all, data provided as grid files.
- **Some caveats:**
 - factorization assumption only sensible **without additional hard jets**;
→ EW corrections to ZZ+jet would have to be included in this configuration.
 - Ansatz **does not include** corrections to non-resonant or off-shell contributions.

Two-by-two kinematics distorted by additional QCD radiation → mapping on two-body phase space needed!

1 \hat{s}, \hat{t} from matching of kinematics:

$$\hat{s} = M_{V_1 V_2}^2; \quad \hat{t} = ?$$

- small p_T of $V_1 V_2 \implies \hat{\Theta} =$ scattering angle in $V_1 V_2$ restframe (collinear or soft QCD corrections)
- large p_T of $V_1 V_2$ (associated with hard jet!)
 define directions of incoming protons in $V_1 V_2$ restframe

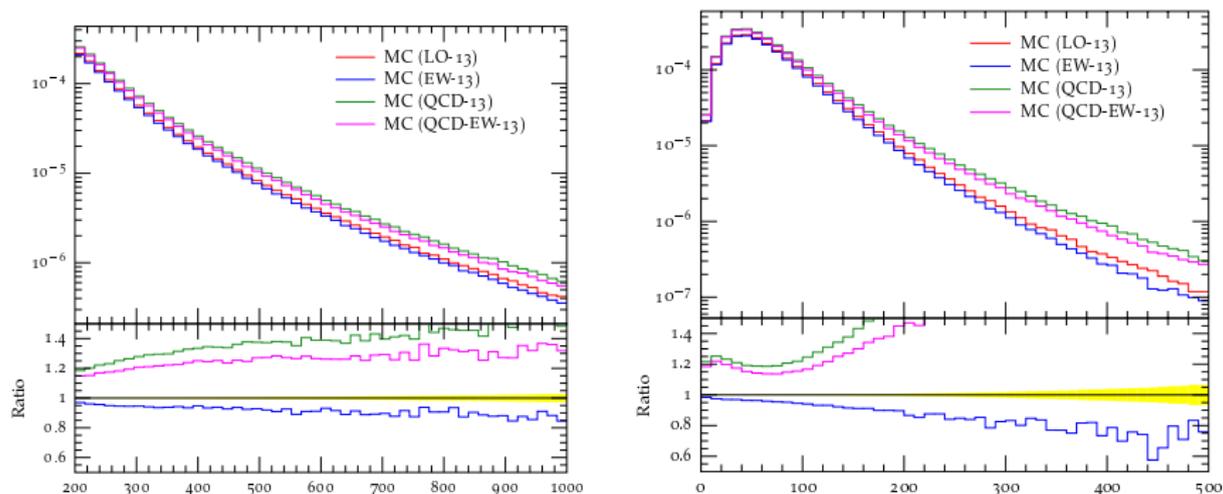
$$\mathbf{e}_{1,2}^* \equiv \frac{\mathbf{p}_{1,2}^*}{|\mathbf{p}_{1,2}^*|}$$

$$\text{scattering axis: } \hat{\mathbf{e}}^* \equiv \frac{\mathbf{e}_1^* - \mathbf{e}_2^*}{|\mathbf{e}_1^* - \mathbf{e}_2^*|}$$

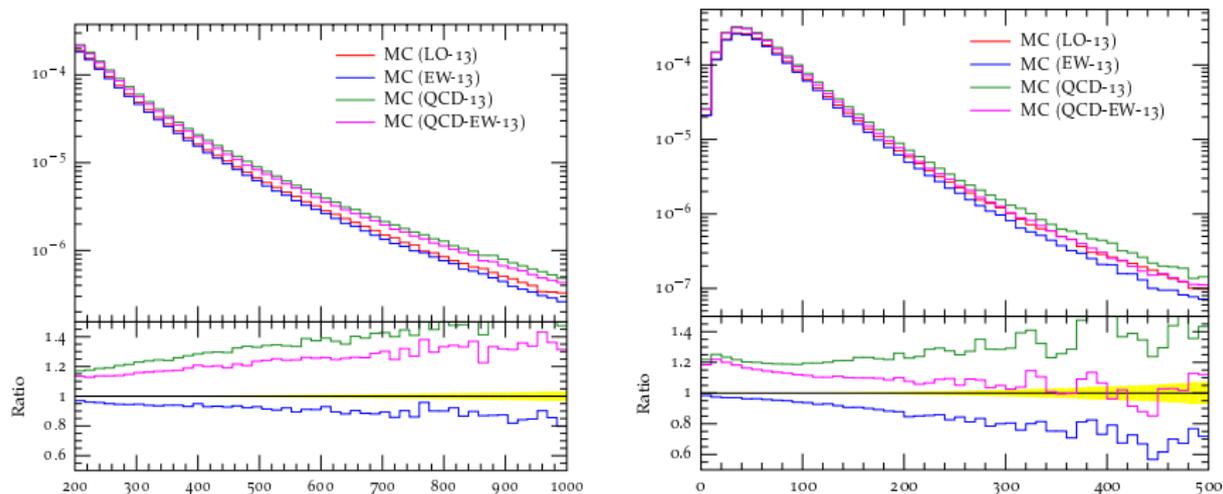
$$\text{scattering angle: } \cos \hat{\Theta} = \frac{\mathbf{p}_{V_1}^*}{|\mathbf{p}_{V_1}^*|} \cdot \hat{\mathbf{e}}^*$$

2 or (HERWIG++)

\hat{s}, \hat{t} provided by the generator

Simulation for $pp \rightarrow ZZ \rightarrow e^+e^-\mu^+\mu^- + X$ at 13 TeV, M_{ZZ} and $p_{T,Z}$ distributions

- Standard Herwig++ setup used (v2.6.3, with simple add-on for EW corrections, 10M events), ZZ at NLO QCD matched with parton showers, hadronization included, underlying event switched off
- huge QCD corrections at large $p_{T,Z}$, factorized ansatz not justified!

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- huge QCD corrections at large $p_{T,Z}$, factorized ansatz not justified!
 \rightarrow jet veto, cut on $p_{T,ZZ}$: $|\sum_i \mathbf{p}_{T,i}| < 0.3 \times \sum_i |\mathbf{p}_{T,i}|$

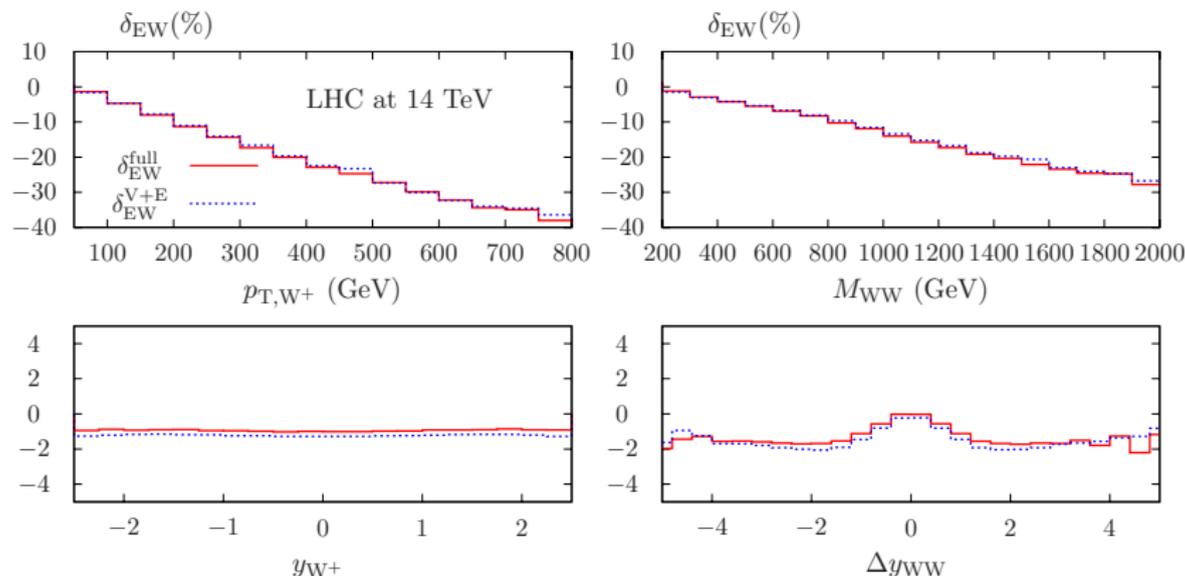
Problem:

- In WW and WZ production **no gauge-invariant separation** of dominant weak corrections and QED possible
- QED contributions inevitably lead to **IR singularities**
- **Real radiation** has to be included:
numerical integration has to deal with singular integrands, check cancellation of divergences, check that slicing cuts drop out, . . .
- Finally, QED effects at the level of 1% (α/π).

Possible solution:

- **V + E approximation**: Endpoint from subtraction contributions \oplus virtual corrections gives IR finite result [Dittmaier 1999]
- Completely avoid computation of real photon radiation

On-shell W-Pair production at LHC14, default cuts

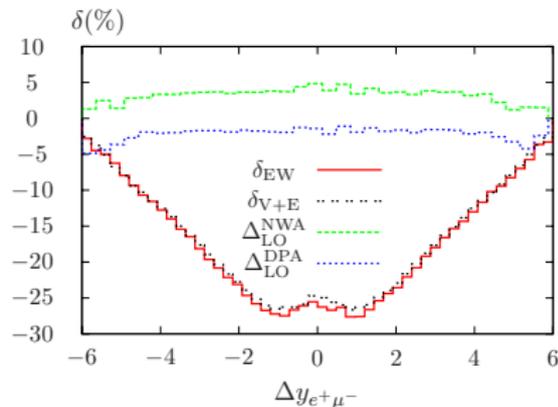
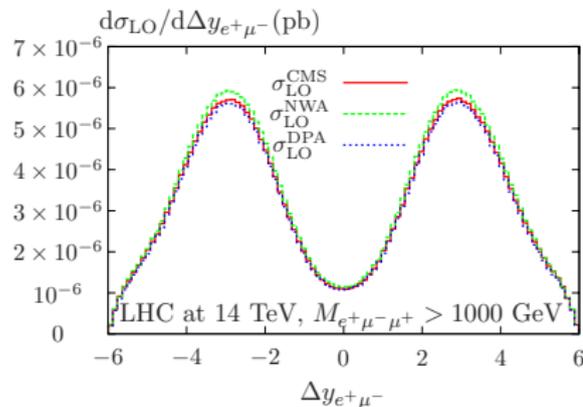


- **Fantastic approximation** of full result (better than 1% in WW, WZ production)
- Approximation works well at high p_T , high invariant masses and near threshold.
- **NNLO EW corrections** at the level of 5–10% at high p_T [Kühn, Metzler, Penin, Uccirati '11]

Conclusion: Corresponding K -factor should be used for MC implementation.

4-Lepton Production – Test of our Approach

$pp \rightarrow (W^+ \rightarrow)e^+ \nu_e (Z \rightarrow)\mu^- \mu^+$ at LHC14, standard event-selection cuts
(Preliminary results!)

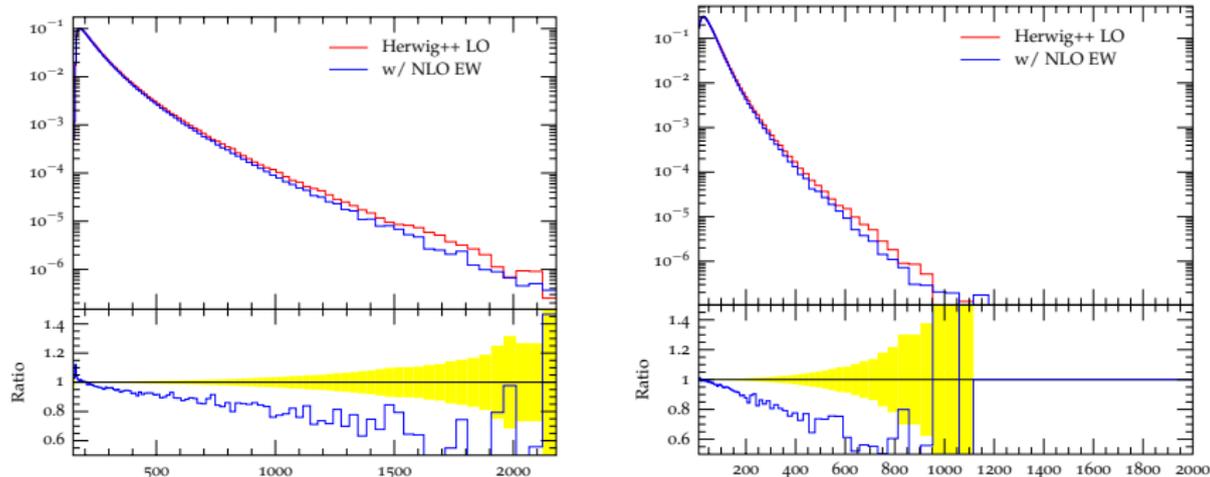


● **LO:** **NWA** and **DPA** work at the level of $\pm 5\%$

● **NLO:**

- δ_{EW} : Full NLO EW corrections to production process in NWA, spin correlations for decay process included
- δ_{V+E} : LO in NWA multiplied with $K_{EW}(\hat{s}, \hat{t})$ (unpolarized K -factor used!)
- **Good agreement** at the 1% level for relative corrections
- Spin correlations well reproduced!

Simulation for $pp \rightarrow (W^+ \rightarrow)e^+\nu_e (W^- \rightarrow)\mu^-\bar{\nu}_\mu + X$ at 8 TeV,
 M_{WW} and $p_{T,W}$ distributions



- Standard Herwig++ setup used (v2.6.2, with simple add-on for EW corrections, 10M events), WW at LO QCD \oplus parton shower, hadronization included, underlying event switched off
- V+E approximate results consistent with [arXiv:1208.3147](https://arxiv.org/abs/1208.3147)

- ✓ **Precise predictions** for vector-boson pair production at NLO QCD & EW exist.
- ✓ **Approximate results** at NNLO available \rightarrow large corrections at high p_T , reduction of residual theoretical uncertainties
- Photon-induced contributions **potentially large** \rightarrow further constrain photon PDFs using LHC data
- ✓ We have computed the full EW corrections to $pp \rightarrow VV$ at hadron colliders
 - **Leptonic decays** have been implemented for WW, ZZ, WZ production, including
 - **spin correlations**
- We have proposed a **straight-forward MC implementation** in the HERWIG++ setup, relying on $2 \rightarrow 2$ K -factors.
 - **Claim:** predictions match the “true” NLO EW result at the level of a few %.
 - QCD uncertainties (PDFs, hadronization, missing higher orders, . . .) presumably much larger
 - Approach could easily be applied to V +jet, $t\bar{t}$ production in the future.

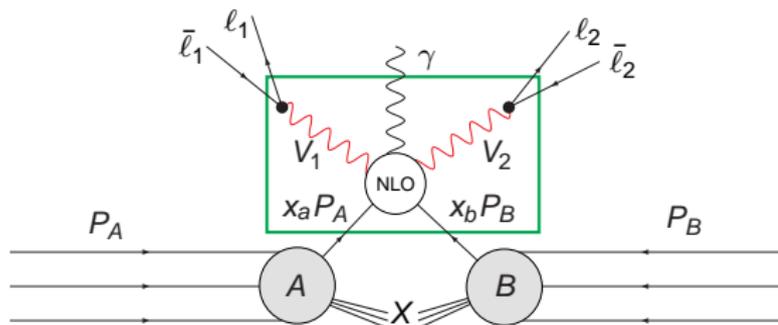
Thank you!

Reminder: Calculation of Hadronic Cross Sections

Schematic illustration for

$$pp \rightarrow V_1 V_2 (+\gamma) + X$$

$$\rightarrow \ell_1 \ell_2 \bar{\ell}_1 \bar{\ell}_2 (+\gamma) + X$$



Hadronic cross sections

$$d\sigma_{AB}(p_A, p_B) = \sum_{a,b} \int_0^1 dx_a \int_0^1 dx_b f_{a/A}(x_a, \mu_F) f_{b/B}(x_b, \mu_F) d\hat{\sigma}_{ab}^{\text{NLO}}(p_a, p_b, \mu_F, \mu_R)$$

$$\times \mathcal{F}^{(4\ell+\gamma)}(\{\mathcal{O}_{\text{FS}}\}), \quad p_{\{a,b\}}^\mu = x_{\{a,b\}} P_{\{A,B\}}^\mu$$

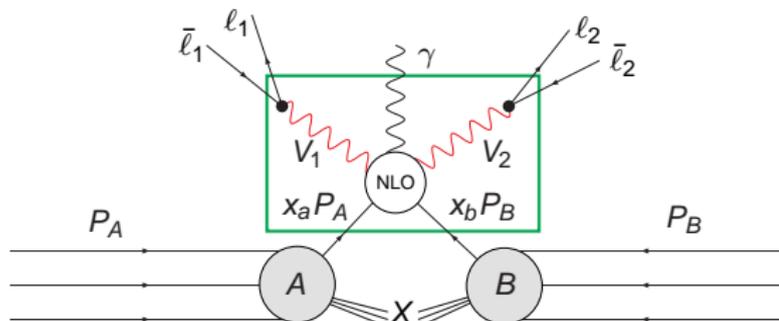
- Dependence on μ_R, μ_F reduced by inclusion of higher perturbative orders
- $\mathcal{F}^{(4\ell+\gamma)}$ incorporates definition of observables + phase-space cuts

Reminder: Calculation of Hadronic Cross Sections

Schematic illustration for

$$pp \rightarrow V_1 V_2 (+\gamma) + X$$

$$\rightarrow \ell_1 \ell_2 \bar{\ell}_1 \bar{\ell}_2 (+\gamma) + X$$



Hadronic cross sections

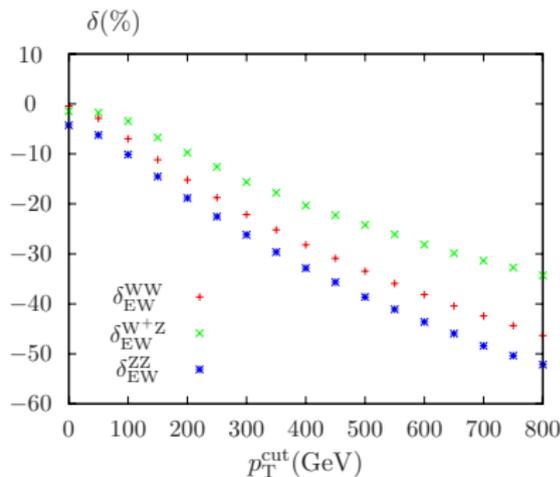
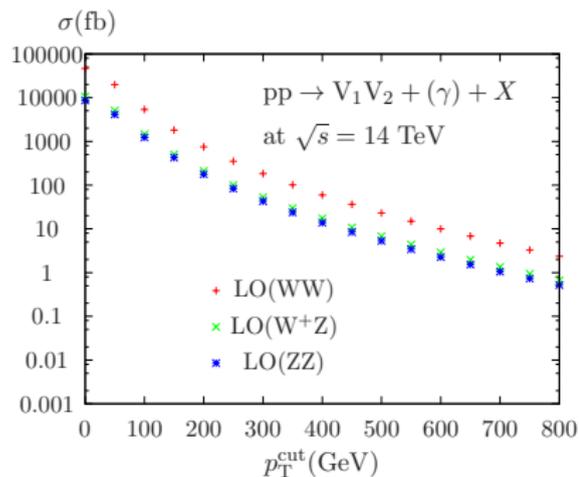
$$d\sigma_{AB}(p_A, p_B) = \sum_{a,b} \int_0^1 dx_a \int_0^1 dx_b f_{a/A}(x_a, \mu_F) f_{b/B}(x_b, \mu_F) d\hat{\sigma}_{ab}^{\text{NLO}}(p_a, p_b, \mu_F, \mu_R)$$

$$\times \mathcal{F}^{(4\ell+\gamma)}(\{\mathcal{O}_{\text{FS}}\}), \quad p_{\{a,b\}}^\mu = x_{\{a,b\}} P_{\{A,B\}}^\mu$$

NLO partonic cross section:

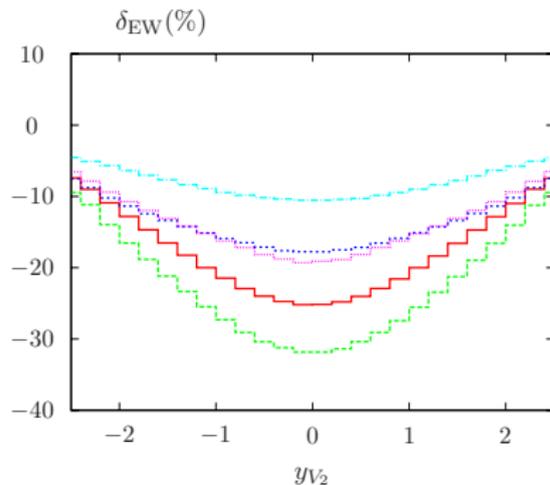
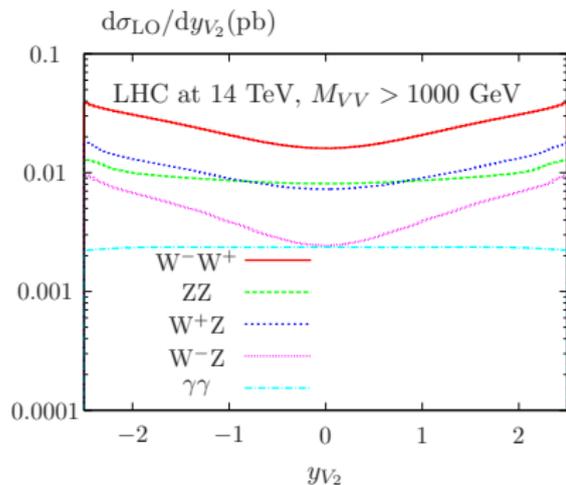
$$\hat{\sigma}_{ab}^{\text{NLO}} = \hat{\sigma}_{ab}^{\text{LO}} + \hat{\sigma}_{ab}^{\text{virt}} + \hat{\sigma}_{ab}^{\text{real}}$$

LHC at 14 TeV, default cuts: $p_{T,V} > 15$ GeV, $|y_V| < 2.5$



- **LO:** Drastically decreasing cross sections for large p_T
- **NLO:**
 - Full result, i.e. virtual, soft, hard, and collinear photons included
 - corrections largest for ZZ, smallest for $\gamma\gamma$

LHC at 14 TeV, high-energy cuts: $p_{T,V} > 15$ GeV, $|y_V| < 2.5$, $M_{VV} > 1000$ GeV



- significant distortion of rapidity distributions at large invariant masses
- Corrections could be misinterpreted as signal of anomalous couplings.

- **Lowest order:** Amplitude given as a product of **on-shell (OS) production amplitude** \otimes **on-shell decay amplitude** \otimes **Breit-Wigner:**

$$\mathcal{M}_{\text{Born,DPA}}^{\bar{q}_1 q_2 \rightarrow V_1 V_2 \rightarrow 4f} = \frac{1}{k_1^2 - M_1^2 + iM_1\Gamma_1} \frac{1}{k_2^2 - M_2^2 + iM_2\Gamma_2} \times \sum_{\lambda_1, \lambda_2} \mathcal{M}_{\text{Born}}^{\bar{q}_1 q_2 \rightarrow V_1, \lambda_1 V_2, \lambda_2} \mathcal{M}_{\text{Born}}^{V_1, \lambda_1 \rightarrow f_3 \bar{f}_4} \mathcal{M}_{\text{Born}}^{V_2, \lambda_2 \rightarrow f_5 \bar{f}_6}$$

- Use **OS-projected momenta** \hat{k} [Denner, Dittmaier, Roth, Wackerath 2000] in the OS matrix elements:

$$\hat{k}_{1,0} = \frac{1}{2}\sqrt{\hat{s}}, \quad \hat{\mathbf{k}}_1 = \frac{\mathbf{k}_1}{|\mathbf{k}_1|} \beta_W \frac{\sqrt{\hat{s}}}{2}, \quad \dots$$

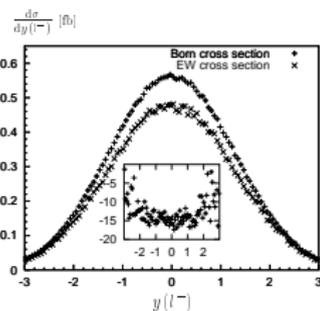
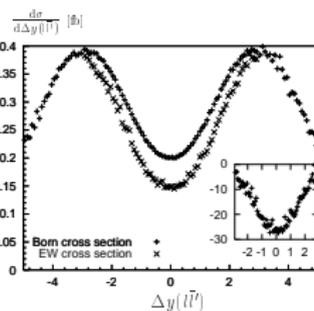
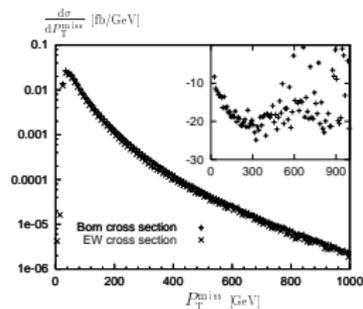
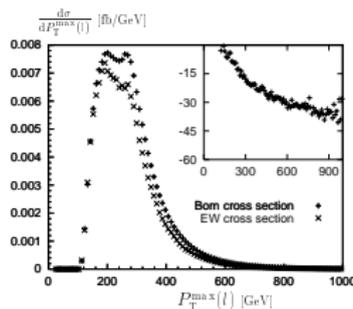
- **NLO:** EW corrections consist of **factorizable** and **non-factorizable** contributions, e.g.

$$\mathcal{M}_{\text{fact}} = \frac{R(k_1, k_2, \theta)}{(k_1^2 - M_1^2 + iM_1\Gamma_1) (k_2^2 - M_2^2 + iM_2\Gamma_2)}$$

Caution: Gauge invariance!

EW corrections to $pp \rightarrow W^+W^- \rightarrow \nu_e e^+ \mu^- \bar{\nu}_\mu$ (DPA)

- Standard LHC event selection cuts applied to final-state leptons and missing transverse momentum; **additionally $M_{e^+\mu^-} > 500$ GeV required**
- Large negative corrections at large transverse momenta
- **Substantial negative corrections** to inclusive observables
- Error due to DPA about 10% in the relative corrections
- **EW corrections significantly larger than experimental error throughout the whole energy range (for $L \sim 30 \text{ fb}^{-1}$)**



[Accomando, Denner, Kaiser: arXiv:0409247 [hep-ph]]

- Simple LL ansatz for $f_{\gamma/p}(x, Q_0^2)$

$$f_{\gamma/p}(x, Q_0^2) = \frac{\alpha}{2\pi} \left[\frac{4}{9} \ln \left(\frac{Q_0^2}{m_u^2} \right) f_{u/p,v}(x, Q_0^2) + \frac{1}{9} \ln \left(\frac{Q_0^2}{m_d^2} \right) f_{d/p,v}(x, Q_0^2) \right] \otimes \frac{1 + (1-x)^2}{x}$$

- Running of $f_{q/p}(x, Q^2)$ at $\mathcal{O}(\alpha)$ affected by photon PDFs!

$$\frac{\partial f_{q/p}(x, \mu^2)}{\partial \ln \mu^2} = \frac{\alpha}{2\pi} \int_x^1 \frac{dy}{y} \left[P_{qq}(y) Q_q^2 f_{q/p}(x/y, \mu^2) + P_{q\gamma}(y) Q_q^2 f_{\gamma/p}(x/y, \mu^2) \right]$$

- Momentum conservation

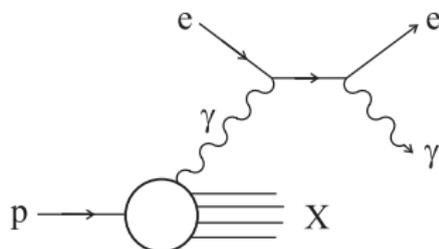
$$\int_0^1 dx \, x \left[\sum_q f_{q/p}(x, \mu^2) + f_{g/p}(x, \mu^2) + f_{\gamma/p}(x, \mu^2) \right] = 1$$

- ⇒ QED effects on $f_{q/p}(x, \mu^2)$ small!
- ⇒ Still large conceptual uncertainties in $f_{\gamma,0}$

Consider the DIS process

$$ep \rightarrow e\gamma + X$$

with high- p_T back-to-back e, γ in the final state



$$\sigma(ep \rightarrow e\gamma + X) = \int dx^\gamma f_{\gamma/p}(x^\gamma, \mu^2) \hat{\sigma}(e\gamma \rightarrow e\gamma),$$

related to Compton scattering

- $x^\gamma = \frac{E_T^\gamma E_e \exp(\eta^\gamma)}{2E_p E_e - E_T^\gamma E_e \exp(-\eta^\gamma)}$
- $f_{\gamma/p}(x^\gamma, \mu^2)$ could be in principle extracted from HERA data!

- $\alpha(0)$: On-shell definition in the Thomson-limit (zero momentum transfer)

$$\bar{u}(p)\Gamma_{\mu}^{Ae\bar{e}}(p,p)u(p)\Big|_{p^2=m_e^2} = e(0)\bar{u}(p)\gamma_{\mu}u(p), \alpha(0) = e(0)^2/4\pi$$

- $\alpha(M_Z)$ obtained via renormalization-group running from 0 to weak scale M_Z

$$\alpha(M_Z) = \frac{\alpha(0)}{1 - \Delta\alpha(M_Z)}, \quad \Delta\alpha(M_Z) = \Pi_{f\neq t}^{AA}(0) - \text{Re} \Pi_{f\neq t}^{AA}(M_Z^2)$$

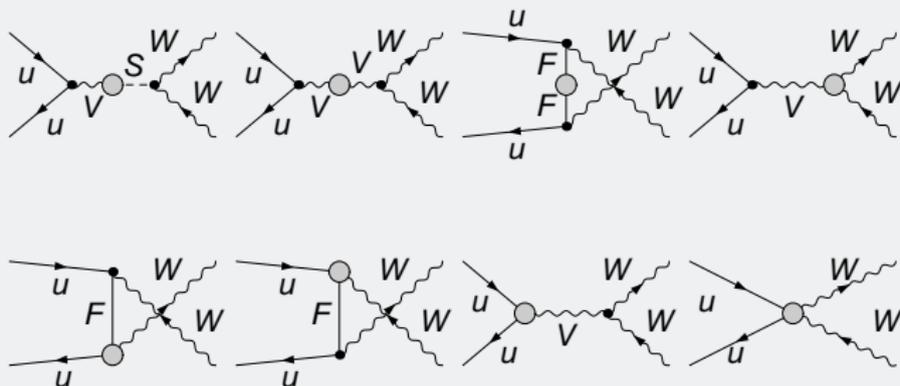
- $\alpha_{G_{\mu}}$ defined through the Fermi constant related to the muon lifetime

$$\alpha_{G_{\mu}} = \frac{\sqrt{2}G_{\mu}M_W^2s_w^2}{\pi} = \frac{\alpha(0)}{1 - \Delta r}$$

Δr includes corrections to muon lifetime not contained in QED-improved Fermi model

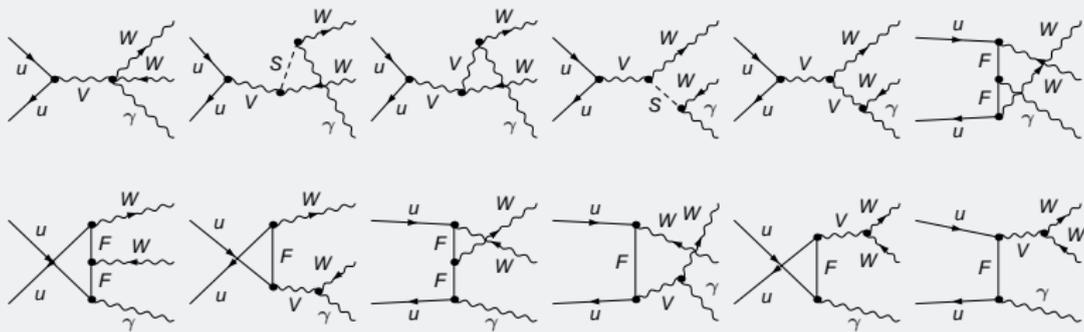
- **light-fermion mass logs contained in $\Pi_{f\neq t}^{AA}(0)$ resummed in effective couplings $\alpha(M_Z)$ and $\alpha_{G_{\mu}}$**

One-loop contributions at $\mathcal{O}(\alpha^3)$



- On-shell renormalization of SM parameters
- We use the Fermi scheme to calculate the loop corrections.
→ universal corrections to Δr absorbed in effective LO coupling
- $V_{ij}^{\text{CKM}} = \delta_{ij}$ within the loops → no renormalization of V_{ij}^{CKM}

Real photon radiation at $\mathcal{O}(\alpha^3)$ (generic diagrams): $q\bar{q} \rightarrow W^-W^+ + \gamma$



- **Soft singularities** due to soft photons
- **Initial-state collinear singularities** due to collinear photon radiation off initial-state quarks \rightarrow renormalization of PDFs
- Introduce small **quark mass m_q** and infinitesimal **photon mass λ** to regularize divergences \rightarrow results exhibit unphysical $\ln m_q$ and $\ln \lambda$ terms

Apply phase-space slicing for numerically-stable evaluation of phase-space integral

Two-cut-off phase-space slicing

- Definition of bremsstrahlung phase space:

$$\sigma_{\text{real}} = \int \text{dPS}(W^-W^+\gamma) |\mathcal{M}^\gamma|^2$$

- Phase-space decomposition:

$$\sigma_{\text{real}} = \sigma_{\text{hard}} + \sigma_{\text{soft}} + \sigma_{\text{coll}}$$

- **Soft limit:** $E_\gamma < \Delta E \ll M_W$

$$\sigma_{\text{soft}}(\Delta E) = -\sigma_{\text{LO}} \left[\frac{e^2}{(2\pi)^3} \int_{|\mathbf{k}_\gamma| < \Delta E} \frac{d^3\mathbf{k}_\gamma}{2\sqrt{\mathbf{k}_\gamma^2 + \lambda^2}} \sum_{ij} \frac{\pm Q_i Q_j (p_i p_j)}{(p_i k_\gamma)(p_j k_\gamma)} \right]$$

- **Collinear limit:** $\theta_{q\gamma} < \Delta\theta \ll 1$, $E_\gamma > \Delta E$

$$\sigma_{\text{coll},q}(\Delta E, \Delta\theta) = \frac{\alpha Q_q^2}{2\pi} \int_0^{1-2\Delta E/\sqrt{s}} dz \frac{(1+z^2)}{1-z} \left(\ln \frac{\hat{s}(\Delta\theta)^2}{4m_q^2} - \frac{2z}{1+z^2} \right) \sigma_{\text{LO}}(z\hat{s})$$

- **Hard bremsstrahlung:** $\theta_{q\gamma} > \Delta\theta$, $E_\gamma > \Delta E$;
numerical evaluation of $\sigma_{\text{hard}}(\Delta E, \Delta\theta)$ without regulators
- Numerical result independent of $\ln \Delta E$ and $\ln \Delta\theta$

$\ln m_q$ and $\ln \lambda$ terms cancel in the sum $\sigma_{\text{virt}} + \sigma_{\text{soft}} + \sigma_{\text{coll}}$ in infrared-safe observables

A problem with unstable particles

Naive implementation of finite width in gauge-boson propagator:

$$\frac{-ig^{\mu\nu}}{q^2 - M_W^2 + i\epsilon} \rightarrow \frac{-ig^{\mu\nu}}{q^2 - M_W^2 + iM_W\Gamma_W}$$

Γ_W includes Dyson summation of self energies, mixing of perturbative orders
→ **might destroy gauge invariance (even at leading order!)**

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Γ_W includes Dyson summation of self energies, mixing of perturbative orders
→ **might destroy gauge invariance (even at leading order!)**

→ **CMS universal solution that**

- respects gauge invariance
- is valid in all phase-space regions

Straightforward implementation:

● **LO:** $M_V^2 \rightarrow \mu_V^2 = M_V^2 - iM_V\Gamma_V$, $\cos^2 \Theta_W = \frac{\mu_W^2}{\mu_Z^2}$, $V = W, Z$

● **NLO:**

- Complex renormalization: $\mathcal{L}_0 \rightarrow \mathcal{L} + \delta\mathcal{L}$, **bare (real) Lagrangian unchanged!**
- Evaluate loop integrals with complex masses

Virtual corrections computed in the `FeynArts/FormCalc/LoopTools (FF)` framework [(FA): Küblbeck, Böhm, Denner 1990; (FC,LT): Hahn, Pérez Victoria 1999; Hahn 2001; (FF): van Oldenborgh, Vermaseren 1990]

1 `FeynArts-3.5`:

- Automatic generation of diagrams
- Calculation of amplitudes

2 `FormCalc-6.1`:

- Algebraical simplification of amplitudes, introduction of tensor coefficients
- Analytical calculation of squared amplitudes
- Spin-, colour- and polarization sums
- Generation of `Fortran` code

3 `LoopTools-2.5`:

- Numerical Passarino–Veltman reduction within `Fortran`
- Numerically-stable evaluation of scalar integrals

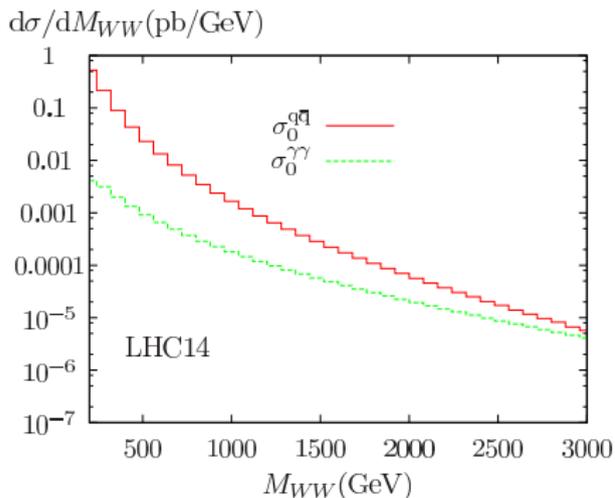
Bremsstrahlung amplitudes computed with `FeynArts/FeynCalc` \oplus `Madgraph` [Alwall et al.], numerical phase-space integration within `Fortran` using the `Vegas` algorithm

pp \rightarrow ZZ + X					
ZZ polarizations	summed	LL	L+	++	+-
LHC8					
$\sigma_{\text{LO}}/\text{pb}$	3.810	0.223	0.396	$10^{-1} \times [0.559]$	2.676
$\delta\sigma_{\text{weak}}/\text{pb}$	-0.179(-0.155)	-0.009(-0.008)	-0.016(-0.014)	-0.002(-0.002)	-0.134(-0.117)
$p_{T,Z} > 500 \text{ GeV}$					
$\sigma_{\text{LO}}/\text{pb}$	$10^{-2} \times [0.101]$	$10^{-7} \times [0.202]$	$10^{-5} \times [0.779]$	$10^{-8} \times [0.504]$	$10^{-3} \times [0.996]$
$\delta\sigma_{\text{weak}}/\text{pb}$	-0.039(-0.030)	-0.975(+0.748)	-0.204(-0.157)	-0.895(-0.425)	-0.383(-0.293)
$p_{T,Z} > 1000 \text{ GeV}$					
$\sigma_{\text{LO}}/\text{pb}$	$10^{-5} \times [0.919]$	$10^{-10} \times [0.121]$	$10^{-7} \times [0.231]$	$10^{-11} \times [0.303]$	$10^{-5} \times [0.915]$
$\delta\sigma_{\text{weak}}/\text{pb}$	-0.557(-0.387)	-2.599(+14.909)	-0.098(-0.070)	-1.742(+1.043)	-0.555(-0.387)

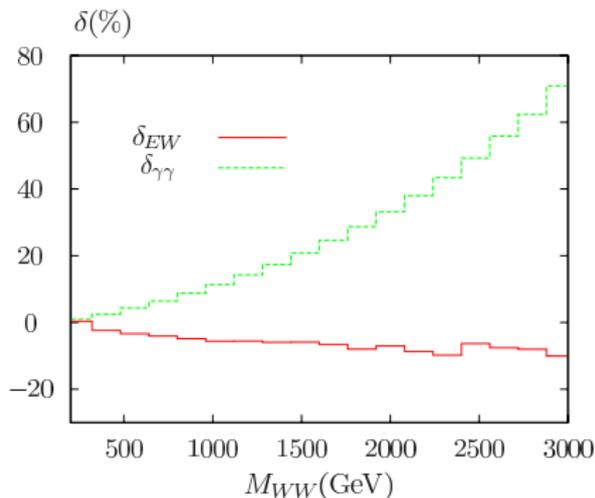
- **small transverse momenta:** 70% from (+-), similar *K*-factors for all polarizations.
- **large transverse momenta:** 99% from (+-), other contributions negligible.
- **Note:** One-loop squared term (given in brackets) contributes at $\sim 10\%$
 → **large uncertainties** due to missing EW higher orders.

Conclusion: One *K*-factor sufficient to describe polarized ZZ production

No cuts

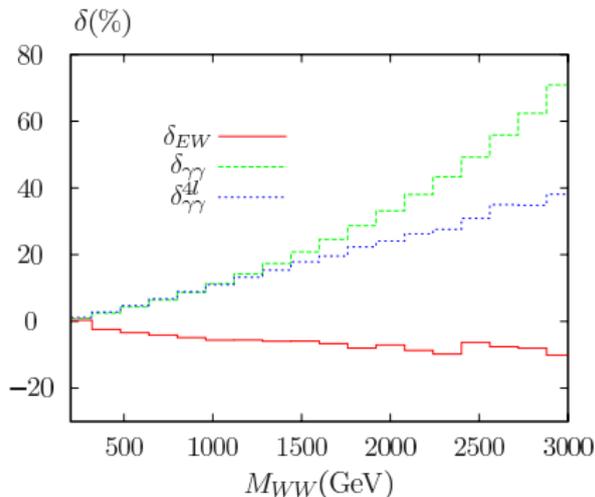
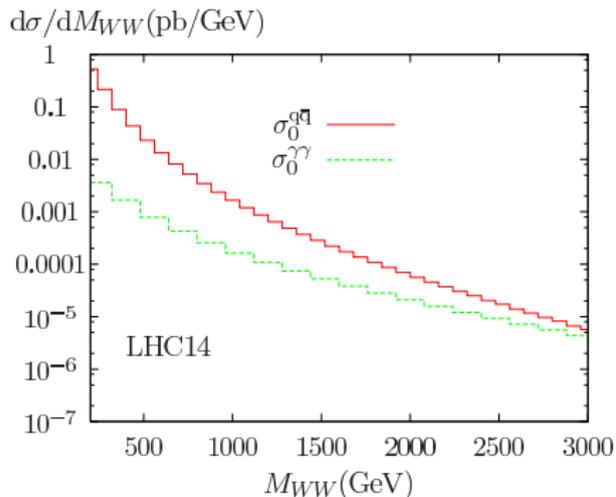


- LO cross section dominated by $q\bar{q}$ contributions
- Rapid decrease of cross section for increasing invariant masses



- EW corrections small even for large values of M_{WW}
 - Large contributions (+80%!) from $\gamma\gamma \rightarrow WW$ at high invariant masses
- ⇒ Leptonic decays?

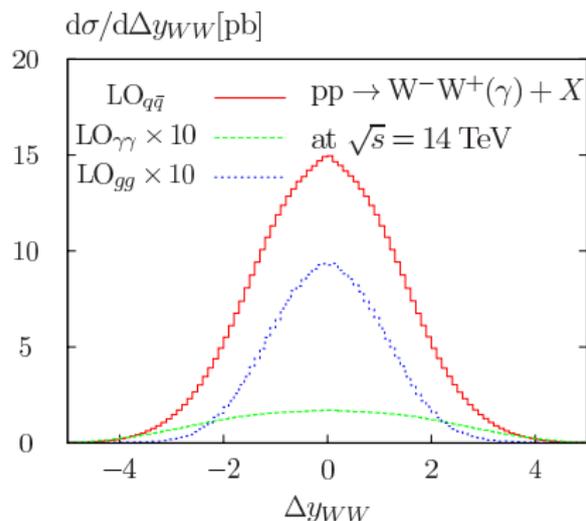
LHC acceptance cuts



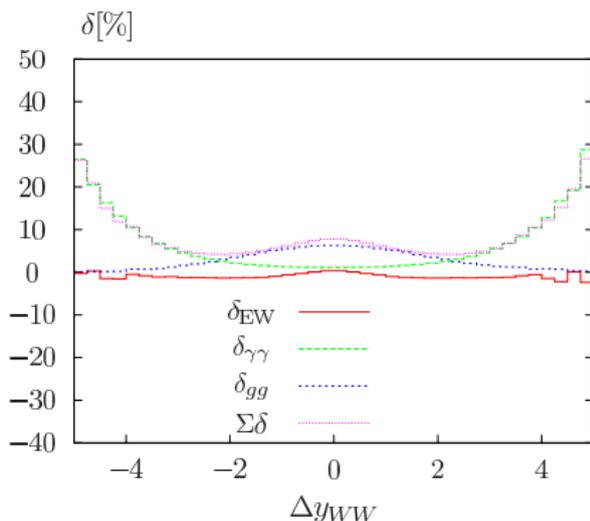
- LO cross section dominated by $q\bar{q}$ contributions
- Rapid decrease of cross section for increasing invariant masses

- Employ LHC cuts on decay products:
 $p_{T,l} > 20$ GeV, $|y_l| < 3$, $p_{T,miss} > 25$ GeV
 \Rightarrow relative effect of $\gamma\gamma \rightarrow WW$ reduced by factor 2 at large M_{WW}

Default cuts: $p_{T,W^\pm} > 15$ GeV, $y_{W^\pm} < 2.5$

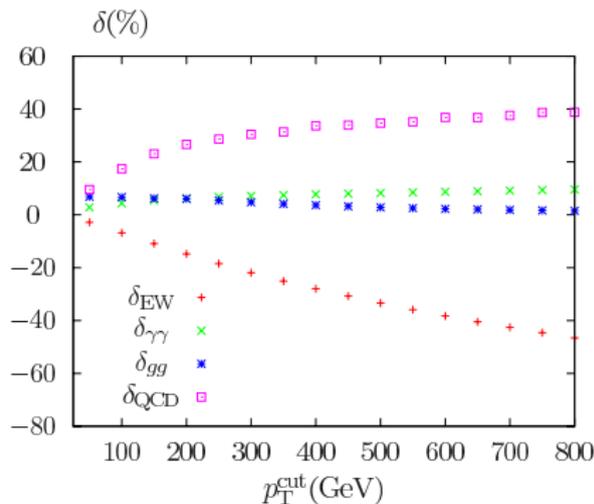
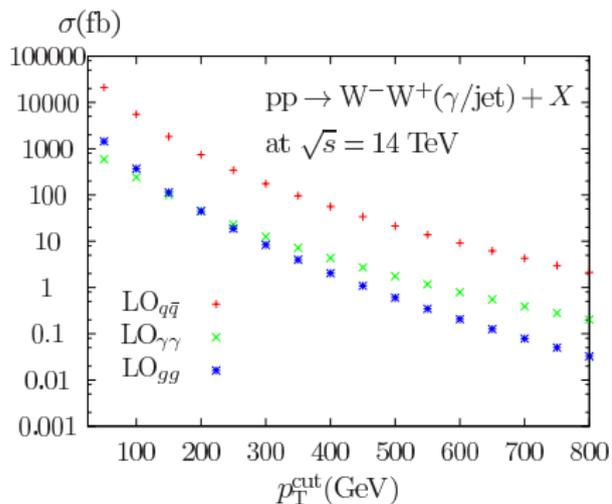


- WW production dominated by events near threshold, isotropic production at small Δy_{WW}
- 5% increase of cross section by gg channel



- EW corrections at the percent level
- Sizable contributions from $\gamma\gamma$ at large $|\Delta y_{WW}|$

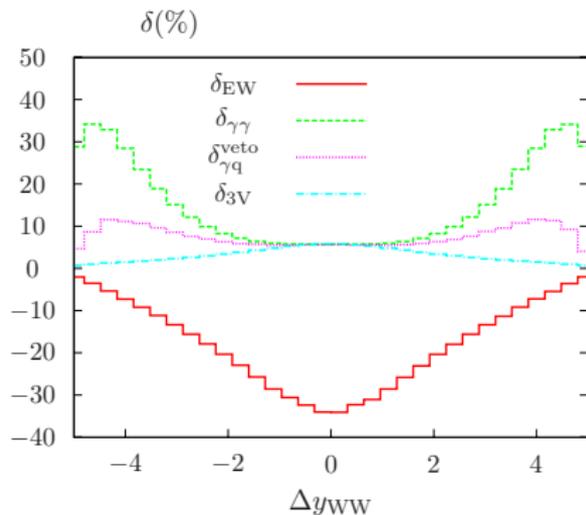
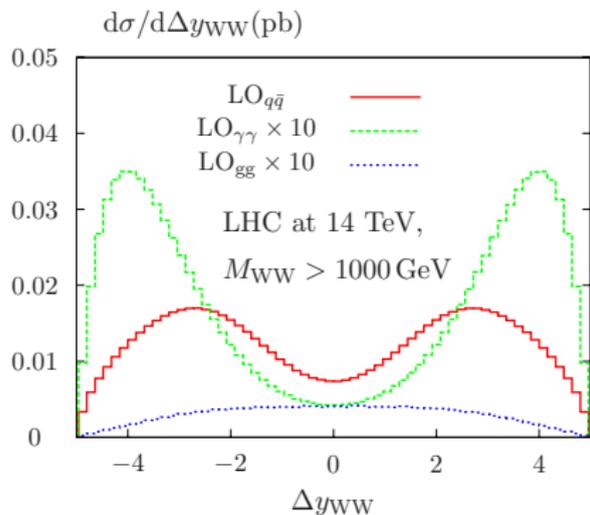
Default cuts: $p_{T,W^\pm} > 15$ GeV, $|y_{W^\pm}| < 2.5$



- assume $\int \mathcal{L} dt = 200 \text{ fb}^{-1}$
 \Rightarrow 1000 WW events with $p_T > 500$ GeV
- decreasing admixture of gg,
 increasing admixture of $\gamma\gamma$

- large admixture of $\gamma\gamma$ (10%!)
- large negative EW corrections (-45%),
 comparable to QCD corrections

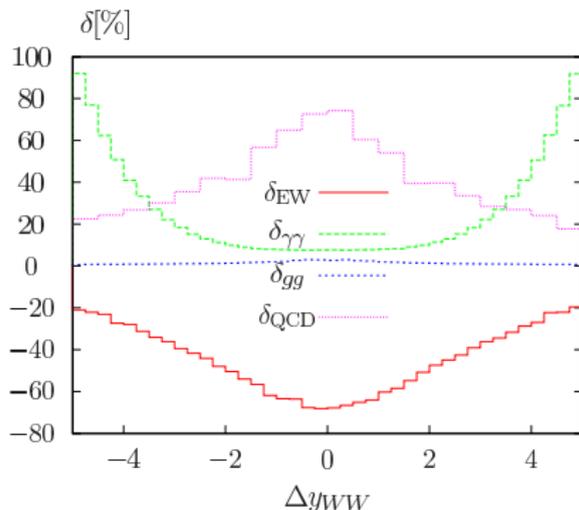
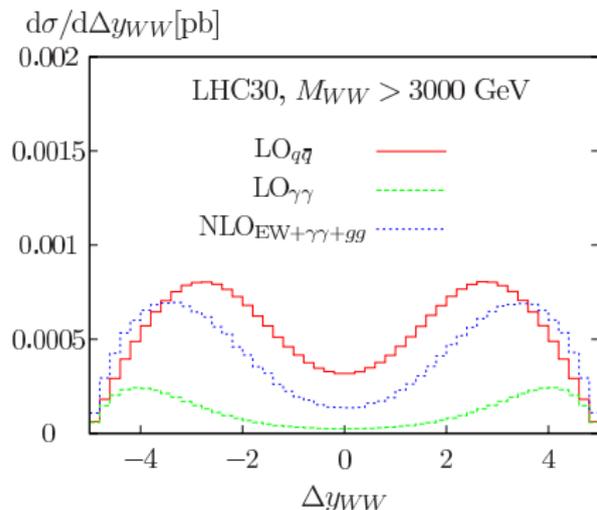
High-energy cuts: $p_{T,W^\pm} > 15 \text{ GeV}$, $y_{W^\pm} < 2.5$, $M_{WW} > 1 \text{ TeV}$



- WW production dominated by small scattering angles
- drastic forward-backward peaking of $\gamma\gamma \rightarrow WW$

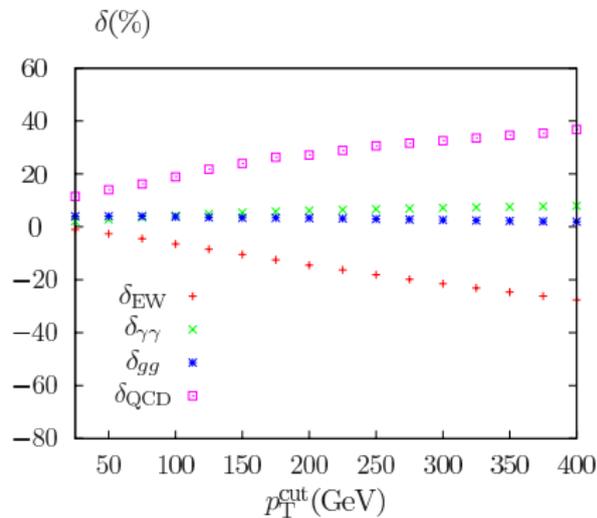
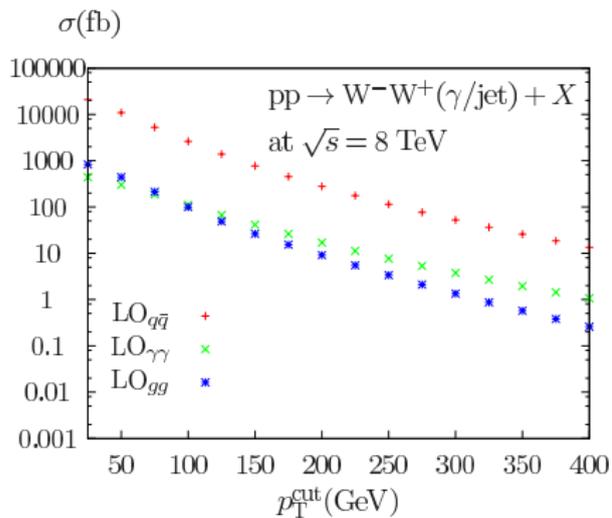
- drastic distortion of angular distribution
- $\Sigma\delta$ varies from -30% and $+45\%$ for $M_{WW} > 1 \text{ TeV}$!

Very-high-energy cuts: $p_{T,W^\pm} > 15$ GeV, $y_{W^\pm} < 2.5$, $M_{WW} > 3$ TeV

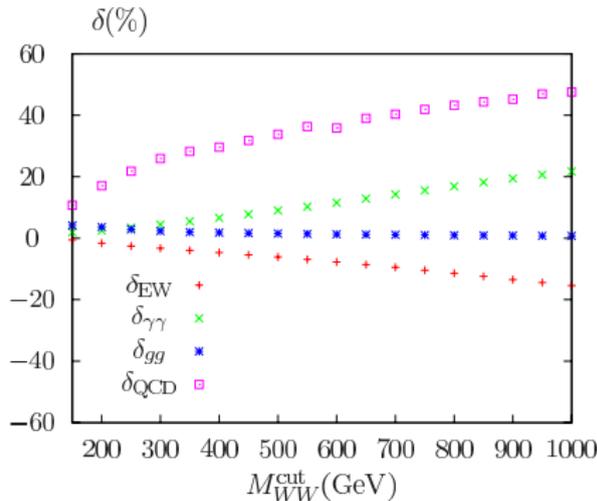
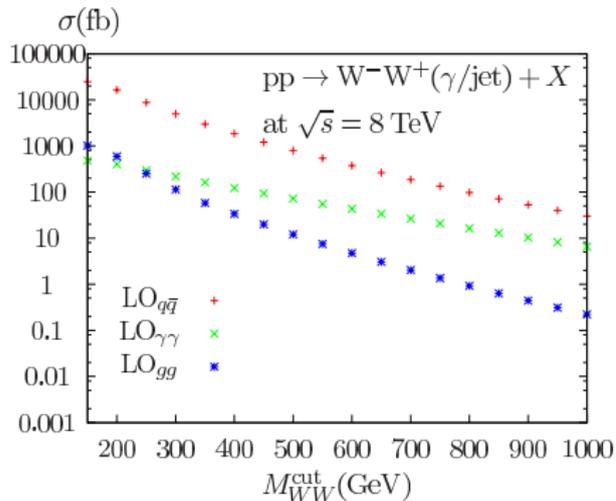


- NLO EW as important as QCD
- extreme distortion due to $\gamma\gamma$ (**caveat:** high uncertainty in photon PDFs)

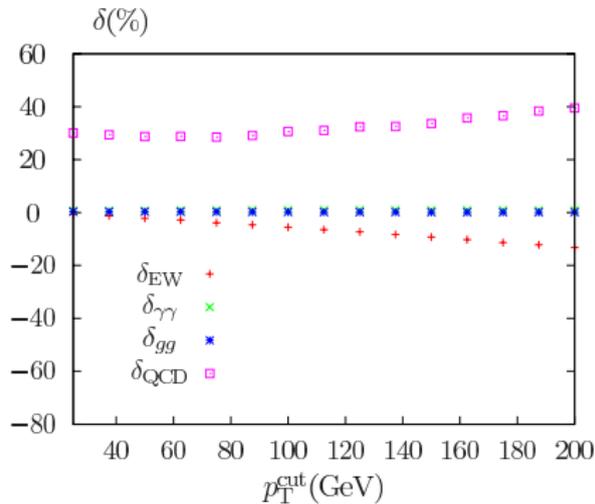
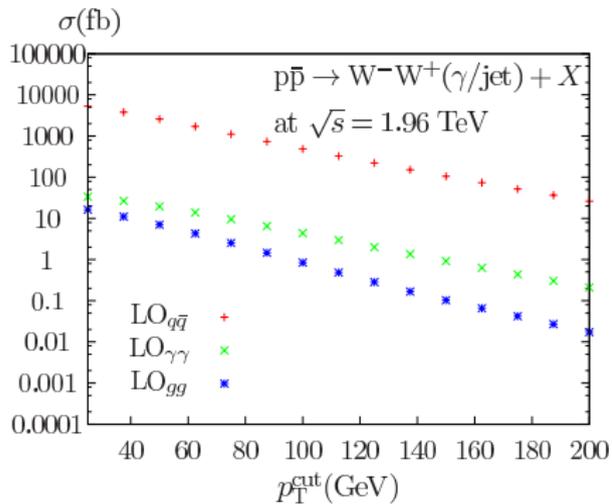
Transverse-momentum distribution at the LHC8



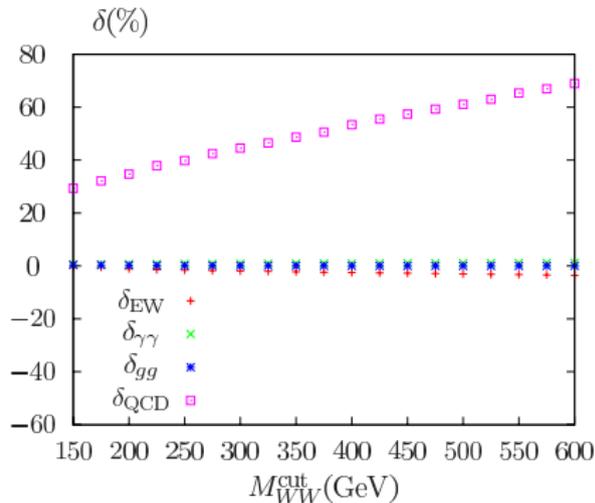
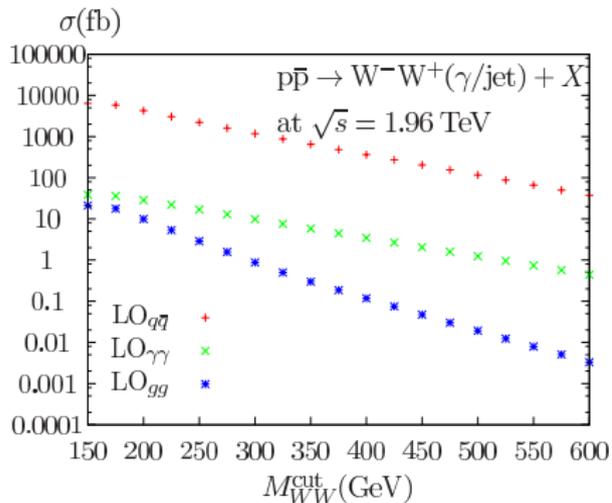
Invariant-mass distribution at the LHC8

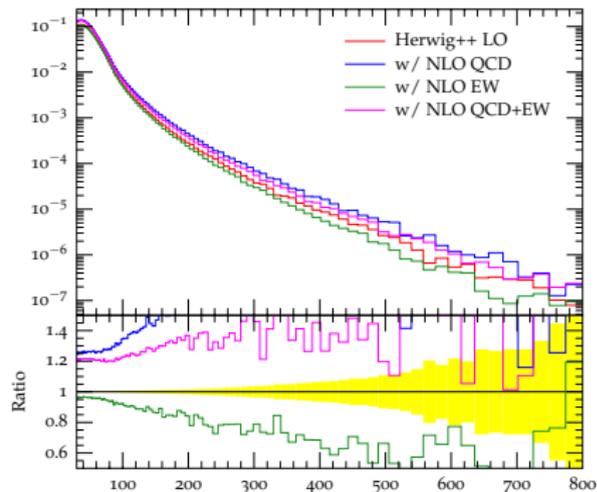
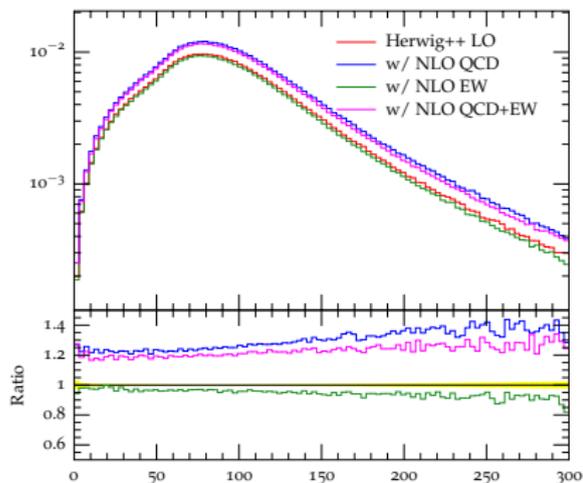


Transverse-momentum distribution at the Tevatron



Invariant-mass distribution at the Tevatron



Simulation for $pp \rightarrow ZZ \rightarrow e^+e^-\mu^+\mu^- + X$ at 8 TeV, $M_{e^+\mu^-}$ and $p_{T,l}$ distributions

- Standard Herwig++ setup used (v2.6.2, with simple add-on for EW corrections, 10M events), ZZ at NLO QCD matched with parton showers, hadronization included, underlying event switched off
- Implementation seems to work fine