## AXEL-2018 <br> Introduction to Particle Accelerators

Transverse optics 2:
$\checkmark$ Hill's equation
$\checkmark$ Phase Space
$\checkmark$ Emittance \& Acceptance
$\checkmark$ Matrix formalism

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## Hill's equation

$\checkmark$ The betatron oscillations exist in both horizontal and vertical planes.
$\checkmark$ The number of betatron oscillations per turn is called the betatron tune and is defined as Qx and Qy .
$\checkmark$ Hill's equation describes this motion mathematically

$$
\frac{d^{2} x}{d s^{2}}+K(s) x=0
$$

$\checkmark$ If the restoring force, $K$ is constant in ' $s$ ' then this is just a Simple Harmonic Motion.
$\checkmark$ ' $s$ ' is the longitudinal displacement around the accelerator.
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## Hill's equation (2)

$\checkmark$ In a real accelerator K varies strongly with 's'.
$\checkmark$ Therefore we need to solve Hill's equation for $K$ varying as a function of ' $s$ '

$$
\frac{d^{2} x}{d s^{2}}+K(s) x=0
$$

$\checkmark$ Remember what we concluded on the mechanical equivalent concerning the shape of the gutter.
$\checkmark$ The phase advance and the amplitude modulation of the oscillation are determined by the shape of the gutter.
$\checkmark$ The overall oscillation amplitude will depend on the initial conditions, i.e. how the motion of the ball started.

## Solution of Hill's equation (1)

$$
\frac{d^{2} x}{d s^{2}}+K(s) x=0
$$

$\checkmark$ Remember, this is a $2^{\text {nd }}$ order differential equation.
$\checkmark$ In order to solve it lets try to guess a solution:

$$
x=\sqrt{\varepsilon \cdot \beta(s)} \cos \left(\phi(s)+\phi_{0}\right)
$$

$\checkmark \varepsilon$ and $\Phi_{0}$ are constants, which depend on the initial conditions.
$\checkmark \beta(s)=$ the amplitude modulation due to the changing focusing strength.
$\checkmark \Phi(s)=$ the phase advance, which also depends on focusing strength.
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## Solution of Hill's equation (2)

$\checkmark$ Define some parameters
$\checkmark \ldots$ and let $\phi=\left(\phi(s)+\phi_{0}\right)$
$x=\sqrt{\varepsilon} . \omega(\mathrm{s}) \cos \phi$
Remember $\Phi$ is still
a function of $s$

$\checkmark$ In order to solve Hill's equation we differentiate our guess, which results in:

$$
x^{\prime}=\sqrt{\varepsilon} \frac{d \omega}{d s} \cos \phi-\sqrt{\varepsilon} \omega \phi^{\prime} \sin \phi
$$

$\checkmark$......and differentiating a second time gives:

$$
x^{\prime \prime}=\sqrt{\varepsilon} \omega^{\prime \prime} \cos \phi-\sqrt{\varepsilon} \omega^{\prime} \phi^{\prime} \sin \phi-\sqrt{\varepsilon} \omega^{\prime} \phi^{\prime} \sin \phi-\sqrt{\varepsilon} \omega \phi^{\prime \prime} \sin \phi-\sqrt{\varepsilon} \omega \phi^{\prime \prime} \cos \phi
$$

$\checkmark$ Now we need to substitute these results in the original equation.

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## Solution of Hill's equation (3)

$\checkmark$ So we need to substitute $x=\sqrt{\varepsilon \cdot \beta(s)} \cos \left(\phi(s)+\phi_{0}\right)$ and its second derivative
$x^{\prime \prime}=\sqrt{\varepsilon} \omega^{\prime \prime} \cos \phi-\sqrt{\varepsilon} \omega^{\prime} \phi^{\prime} \sin \phi-\sqrt{\varepsilon} \omega^{\prime} \phi^{\prime} \sin \phi-\sqrt{\varepsilon} \omega \phi^{\prime \prime} \sin \phi-\sqrt{\varepsilon} \omega \phi^{\prime \prime} \cos \phi$
into our initial differential equation

$$
\frac{d^{2} x}{d s^{2}}+K(s) x=0
$$

$\checkmark$ This gives:
$\sqrt{\varepsilon} \omega^{\prime \prime} \cos \phi-\sqrt{\varepsilon} \omega^{\prime} \phi^{\prime} \sin \phi-\sqrt{\varepsilon} \omega^{\prime} \phi^{\prime} \sin \phi-\sqrt{\varepsilon} \omega \phi^{\prime \prime} \sin \phi-\sqrt{\varepsilon} \omega \phi^{\prime \prime} \cos \phi$

$$
+K(s) \sqrt{\varepsilon} \omega \cos \phi=0
$$

Sine and Cosine are orthogonal and will never be 0 at the same time
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The sum of the coefficients must vanish separately to make our guess valid for all phases

## Solution of Hill's equation (4)

$$
\begin{gathered}
\sqrt{\varepsilon} \omega^{\prime \prime} \cos \phi-\sqrt{\varepsilon} \omega^{\prime} \phi^{\prime} \sin \phi-\sqrt{\varepsilon} \omega^{\prime} \phi^{\prime} \sin \phi-\sqrt{\varepsilon} \omega \phi^{\prime \prime} \sin \phi-\sqrt{\varepsilon} \omega \phi^{\prime 2} \cos \phi \\
+K(s) \sqrt{\varepsilon} \omega \cos \phi=0
\end{gathered}
$$

$\checkmark$ Using the 'Sin'

$$
\longrightarrow 2 \omega^{\prime} \phi^{\prime}+\omega \phi^{\prime \prime}=0 \longrightarrow 2 \omega \omega^{\prime} \phi^{\prime}+\omega^{2} \phi^{\prime \prime}=0
$$ terms

$\checkmark$ We defined $\beta=\omega^{2}$, which after differentiating gives $\beta^{\prime}=2 \omega \omega^{\prime}$
$\checkmark$ Combining $2 \omega \omega^{\prime} \phi^{\prime}+\omega^{2} \phi^{\prime \prime}=0$ and $\beta^{\prime}=2 \omega \omega^{\prime}$ gives: $\beta^{\prime} \phi^{\prime}+\beta \phi^{\prime \prime}=\left(\beta \phi^{\prime}\right)^{\prime}=0$

As condition for our guessed solution to be valid we get:

$$
\beta \phi^{\prime}=\text { const. }=1 \text { hence } \phi^{\prime}=\frac{d \phi}{d s}=\frac{1}{\beta}
$$

$\checkmark$ So our guess seems to be correct

## Solution of Hill's equation (5)

$\checkmark$ Since our solution was correct we have the following for $x$ :

$$
x=\sqrt{\varepsilon \cdot \beta} \cos \phi
$$

$\checkmark$ For $x$ ' we have now:

$\checkmark$ Thus the expression for $x$ ' finally becomes:

$$
x^{\prime}=-\alpha \sqrt{\varepsilon / \beta} \cos \phi-\sqrt{\varepsilon / \beta} \sin \phi
$$

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## Phase Space Ellipse

$\checkmark$ So now we have an expression for $x$ and $x$ '

$$
x=\sqrt{\varepsilon . \beta} \cos \phi \quad \text { and } \quad x^{\prime}=-\alpha \sqrt{\varepsilon / \beta} \cos \phi-\sqrt{\varepsilon / \beta} \sin \phi
$$

$\checkmark$ If we plot $x^{\prime}$ versus $x$ as $\Phi$ goes from 0 to $2 \pi$ we get an ellipse, which is called the phase space ellipse.


## Phase Space Ellipse (2)

$\checkmark$ As we move around the machine the shape of the ellipse will change as $\beta$ changes under the influence of the quadrupoles
$\checkmark$ However the area of the ellipse ( $\pi \varepsilon$ does not change

$\checkmark \underline{\varepsilon}$ is called the transverse emittance and is determined by the initial beam conditions.
$\checkmark$ The units are meter-radians, but in practice we use more often $\mathrm{mm} \cdot \mathrm{mrad}$.

## Phase Space Ellipse (3)


$\checkmark$ For each point along the machine the ellipse has a particular orientation, but the area remains the same

## Phase Space Ellipse (4)


$\checkmark$ The projection of the ellipse on the $x$-axis gives the Physical transverse beam size.
$\checkmark$ Therefore the variation of $\beta(s)$ around the machine will tell us how the transverse beam size will vary.


## Emittance \& Acceptance

$\checkmark$ To be rigorous we should define the emittance slightly differently.
$\checkmark$ Observe all the particles at a single position on one turn and measure both their position and angle.
$\checkmark$ This will give a large number of points in our phase space plot, each point representing a particle with its co-ordinates $x, x^{\prime}$.

$\checkmark$ The emittance is the area of the ellipse, which contains all, or a defined percentage, of the particles.
$\checkmark$ The acceptance is the maximum area of the ellipse, which the emittance can attain without losing particles.

## Emittance measurement



## Matrix Formalism

$\checkmark$ Lets represent the particles transverse position and angle by a column matrix.

$$
\binom{x}{x^{\prime}}
$$

$\checkmark$ As the particle moves around the machine the values for $x$ and $x$ ' will vary under influence of the dipoles, quadrupoles and drift spaces.
$\checkmark$ These modifications due to the different types of magnets can be expressed by a Transport Matrix M
$\checkmark$ If we know $x_{1}$ and $x_{1}^{\prime}$ at some point $s_{1}$ then we can calculate its position and angle after the next magnet at position $s_{2}$ using:

$$
\binom{x\left(s_{2}\right)}{x\left(s_{2}\right)^{\prime}}=M\binom{x\left(s_{1}\right)}{x\left(s_{1}\right)^{\prime}}=\left(\begin{array}{ll}
a & b \\
c & d
\end{array}\right)\binom{x\left(s_{1}\right)}{x\left(s_{1}\right)^{\prime}}
$$

## How to apply the formalism

$\checkmark$ If we want to know how a particle behaves in our machine as it moves around using the matrix formalism, we need to:
$\checkmark$ Split our machine into separate elements as dipoles, focusing and defocusing quadrupoles, and drift spaces.
$\checkmark$ Find the matrices for all of these components
$\checkmark$ Multiply them all together
$\checkmark$ Calculate what happens to an individual particle as it makes one or more turns around the machine

## Matrix for a drift space

$\checkmark$ A drift space contains no magnetic field.
$\checkmark$ A drift space has length L.


## Matrix for a quadrupole

$\checkmark$ A quadrupole of length $L$.


Remember $\mathrm{B}_{\mathrm{y}} \propto \mathrm{x}$ and the deflection due to the magnetic field is: $\frac{L B_{y}}{(B \rho)}=-\frac{L K}{(B \rho)} \cdot x$


## Matrix for a quadrupole (2)

$\checkmark$ We found:

$$
\binom{x_{2}}{x_{2}{ }^{\prime}}=\left(\begin{array}{cc}
1 & 0 \\
-\frac{L K}{(B \rho)} & 1
\end{array}\right)\binom{x_{1}}{x_{1}^{\prime}}
$$

$\checkmark$ Define the focal length of the quadrupole as $f=\frac{(B \rho)}{K L}$

$$
\binom{x_{2}}{x_{2}^{\prime}}=\left(\begin{array}{cc}
1 & 0 \\
-\frac{1}{f} & 1
\end{array}\right)\binom{x_{1}}{x_{1}^{\prime}}
$$

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## How now further?

$\checkmark$ For our purpose we will treat dipoles as simple drift spaces as they bend all the particles by the same amount.
$\checkmark$ We have Transport Matrices corresponding to drift spaces and quadrupoles.
$\checkmark$ These matrices describe the real discrete focusing of our quadrupoles.
$\checkmark$ Now we must combine these matrices with our solution to Hill's equation, since they describe the same motion......

## Questions....,Remarks...?



